

## Conceptual Questions

7.1. If you were to throw the rocks in the opposite direction you wanted to go, you would be pushed by the rocks in the right direction. Throwing the rocks requires a force to accelerate them (Newton's second law). So you exert a force on the rock in one direction and the rock exerts an equal force on you in the opposite direction (Newton's third law). This force will cause you to slide along the ice in the opposite direction that you threw the rock. Note that you will move most efficiently when you use a horizontal force, which means that you throw the rock horizontally.

7.2. The paddle, you, and the canoe can be treated as a single object. You can push backward on the water with the paddle so that the water pushes forward on the paddle. The figure shows how the backward force of the paddle on the water and the forward force of the water on the paddle are action/reaction pairs. Since you hold the paddle while sitting in the canoe, the force of the water on the paddle causes the paddle-person-canoe object to move forward. The vertical force between the canoe and water we label as a normal force here but will later identify as the buoyant force.

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7.3. The rocket forces the exhaust gases down, and the hot gases push up on the rocket. The two forces are a Newton's third law pair. The rocket accelerates upward because the force of the exhaust gases on the rocket is greater than the force of gravity.

## Sketch Interaction diagram Free-body diagrams


7.4. The player pushes down on the floor, which pushes back up on him. The player accelerates upward because the force of his push is greater than the force of gravity.


Interaction diagram

$\mathrm{P}=$ Player
S = Surface
EE = Entire Earth

## Free-body diagrams


7.5. Newton's third law tells us that the force of the mosquito on the car has the same magnitude as the force of the car on the mosquito.
7.6. The mosquito has a much smaller mass than the car, so the magnitude of the interaction force between the car and mosquito, although equal on each, causes the mosquito to have a much larger acceleration. In fact, the acceleration is usually fatal to the mosquito.
7.7. Newton's third law tells us that the magnitude of the forces are equal. The acceleration of the truck and car are determined by the net force on each.
7.8. The force of the wagon on the girl acts on the girl, whereas the force of the girl on the wagon acts on the wagon. The wagon's motion is determined by the net force acting on it, so if the girl pulls hard enough to overcome any other opposing forces acting on the wagon, the wagon will move forward. So try saying, "But, my dear, the net force on the wagon determines if it will move forward. The forces you mention act on different objects, and so cannot cancel."
7.9. The net force on each team determines that team's motion. The net horizontal force on each team is the difference between the rope's pull and friction with the ground. So the team that wins the tug-of-war is not the team that pulls harder, but the team that is best able to keep from sliding along the ground.
7.10. This technique will not work because the magnet is part of the cart, not external to it. The forces between the magnet and cart have the same magnitude but act in the opposite directions. Therefore, although the two objects may accelerate toward each other, the cart-magnet system as a whole will not move. (Actually, it would be more precise to say that the center of mass of the cart-magnet system does not move.)
7.11. The scale reads 5 kg . The left-hand mass performs a function no different than the ceiling would if the rope were attached to the ceiling (i.e., both pull upward with the force required to suspend 5 kg ). The force of gravity acting on the right-hand mass provides the downward force on the spring scale that the spring scale converts to mass in its display.
7.12. The scale reads 5 kg . The left-hand mass performs a function no different than the wall would if the rope were attached to the wall (i.e., both pull leftward with the force required to suspend 5 kg ). The right-hand mass provides the rightward force on the spring scale that the spring scale converts to mass and displays.

### 7.13.



The figure shows the horizontal forces on blocks B and A using the massless-string approximation in the absence of friction. The hand must accelerate both blocks A and B, so more force is required to accelerate the greater mass. Thus the force of the string on B is smaller than the force of the hand on A .

### 7.14.



The pulley will not rotate. As shown in the free-body diagrams above, the force of gravity pulls down equally on both blocks so the tension forces, which act as if they were a Newton third law pair, pull up equally on each with the same magnitude force as the force of gravity. The net force on each block is therefore zero, so they do not move and the pulley does not rotate.
7.15. Block A's acceleration is greater in case $b$. In case $a$, the hanging 10 N must accelerate both the mass of $A$ and its own mass, leading to a smaller acceleration than case b , where the entire 10 N force accelerates the mass of block A .

| Case a | Case b |
| :---: | :---: |
| $10 \mathrm{~N}=\left(M_{\mathrm{A}}+M_{10 \mathrm{~N}}\right) a$ | $10 \mathrm{~N}=M_{\mathrm{A}} a$ |
| $a=\frac{10 \mathrm{~N}}{\left(M_{\mathrm{A}}+M_{10 \mathrm{~N}}\right)}$ | $a=\frac{10 \mathrm{~N}}{M_{\mathrm{A}}}$ |

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## Exercises and Problems

## Section 7.2 Analyzing Interacting Objects

### 7.1. Visualize:



Solve: (a) The weight lifter is holding the barbell in dynamic equilibrium as he stands up, so the net force on the barbell and on the weight lifter must be zero. The barbells have an upward contact force from the weight lifter and the gravitational force downward. The weight lifter has a downward contact force from the barbells and an upward one from the surface. Gravity also acts on the weight lifter.

$\mathrm{BB}=$ Barbells
WL $=$ Weight lifter
S = Surface EE = Entire Earth
(b) The system is the weight lifter and barbell, as indicated in the figure.
(c)

## Free-body diagrams



### 7.2. Visualize:

## Sketch



Earth (E)

Solve: (a) Both the bowling ball and the soccer ball have a normal force from the surface and gravitational force on them. The interaction forces between the two are equal and opposite.

## Interaction diagram


$\mathrm{SB}=$ Soccer ball
$\mathrm{BB}=$ Bowling ball
$S=$ Surface
EE = Entire Earth
(b) The system consists of the soccer ball and bowling ball, as indicated in the figure.
(c)


Assess: Even though the soccer ball bounces back more than the bowling ball, the forces that each exerts on the other are part of an action/reaction pair, and therefore have equal magnitudes. Each ball's acceleration due to the forces on it is determined by Newton's second law, $a=F_{\text {net }} / m$, which depends on the mass. Since the masses of the balls are different, their accelerations are different.

### 7.3. Visualize:



Solve: (a) Both the mountain climber and bag of supplies have a normal force from the surface on them, as well as a gravitational force vertically downward. The rope has gravity acting on it, along with pulls on each end from the mountain climber and supply bag. The mountain climber experiences static friction with the surface, whereas the bag experiences kinetic friction with the surface.

$\mathrm{MC}=$ Mountain climber
$\mathrm{R}=$ Rope $\mathrm{Su}=$ Supply bag
S = Surface
EE = Entire Earth
(b) The system consists of the mountain climber, rope, and bag of supplies, as indicated in the figure.
(c)

Free-body diagrams


Mountain climber
Rope
Supply bag

Assess: Since the motion is along the surface, it is convenient to choose the $x$-coordinate axis along the surface. The free-body diagram of the rope shows pulls that are slightly off the $x$-axis since the rope is not massless.

### 7.4. Visualize:



Solve: (a) The car and rabbit both experience a normal force and friction from the floor and a gravitational force from the Earth. The push that each exerts on the other is a Newton's third law force pair.

(b) The system consists of the car and stuffed rabbit, as indicated in the figure.
(c)

7.5. Visualize: Please refer to Figure EX7.5.

Solve: (a) Gravity acts on both blocks. Block A is in contact with the floor and experiences a normal force and friction. The string tension is the same on both blocks since the rope and pulley are massless and the pulley is frictionless. There are two third law pair of forces at the surface where the two blocks touch. Block B pushes against Block A with a normal force, while Block A pushes back against Block B. There is also friction between the two blocks at the surface.

Interaction diagram

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(b) A string that will not stretch constrains the two blocks to accelerate at the same rate but in opposite directions. Block A accelerates down the incline with an acceleration equal in magnitude to the acceleration of Block B up the incline. The system consists of the two blocks, as indicated in the figure above.
(c)

## Free-body diagrams



Assess: The inclined coordinate systems allows the acceleration $a$ to be purely along the $x$-axis. This is convenient because it simplifyies the mathematical expression of Newton's second law.
7.6. Visualize: Please refer to Figure EX7.6.

Solve: (a) For each block, there is a gravitational force due to the Earth, a normal force and kinetic friction due to the surface, and a tension force due to the rope.

(b) The tension in the massless ropes over the frictionless pulley is the same on both blocks. Block A accelerates down the incline with the same magnitude acceleration that Block B has up the incline. The system consists of the two blocks, as indicated in the figure.
(c)


Assess: The inclined coordinate systems allow the acceleration $a$ to be purely along the $x$-axis. This is convenient since then one component of $a$ is zero, simplifying the mathematical expression of Newton's second law.

## Section 7.3 Newton's Third Law

7.7. Model: We will model the astronaut and the chair as particles. The astronaut and the chair will be denoted by A and C, respectively, and they are separate systems. The launch pad is a part of the environment.

## Visualize:

Pictorial representation




Solve: (a) Newton's second law for the astronaut is

$$
\sum\left(F_{\text {on } \mathrm{A}}\right)_{y}=n_{\mathrm{C} \text { on } \mathrm{A}}-\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=m_{\mathrm{A}} a_{\mathrm{A}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{C} \text { on } \mathrm{A}}=\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=m_{\mathrm{A}} g
$$

By Newton's third law, the astronaut's force on the chair is

$$
n_{\mathrm{A} \text { on } \mathrm{C}}=n_{\mathrm{C} \text { on } \mathrm{A}}=m_{\mathrm{A}} g=(80 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=7.8 \times 10^{2} \mathrm{~N}
$$

(b) Newton's second law for the astronaut is:

$$
\sum\left(F_{\text {on A }}\right)_{y}=n_{\mathrm{C} \text { on } \mathrm{A}}-\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=m_{\mathrm{A}} a_{\mathrm{A}} \Rightarrow n_{\mathrm{C} \text { on } \mathrm{A}}=\left(F_{\mathrm{G}}\right)_{\mathrm{A}}+m_{\mathrm{A}} a_{\mathrm{A}}=m_{\mathrm{A}}\left(g+a_{\mathrm{A}}\right)
$$

By Newton's third law, the astronaut's force on the chair is

$$
n_{\mathrm{A} \text { on } \mathrm{C}}=n_{\mathrm{C} \text { on A }}=m_{\mathrm{A}}\left(g+a_{\mathrm{A}}\right)=(80 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}+10 \mathrm{~m} / \mathrm{s}^{2}\right)=1.6 \times 10^{3} \mathrm{~N}
$$

Assess: This is a reasonable value because the astronaut's acceleration is greater than $g$.
7.8. Visualize: Please refer to Figure EX7.8.


Solve: Since the ropes are massless we can treat the tension force they transmit as a Newton's third law force pair on the blocks. The connection shown in Figure EX7.8 has the same effect as a frictionless pulley on these massless ropes. The blocks are in equilibrium as the mass of A is increased until block B slides, which occurs when the static friction on B is at its maximum value. Applying Newton's first law to the vertical forces on block B gives $n_{\mathrm{B}}=\left(F_{\mathrm{G}}\right)_{\mathrm{B}}=m_{\mathrm{B}} g$. The static friction force on B is thus

$$
\left(f_{\mathrm{s}}\right)_{\mathrm{B}}=\mu_{\mathrm{s}} n_{\mathrm{B}}=\mu_{\mathrm{s}} m_{\mathrm{B}} g .
$$

Applying Newton's first law to the horizontal forces on B gives $\left(f_{\mathrm{s}}\right)_{\mathrm{B}}=T_{\mathrm{A} \text { on } \mathrm{B}}$, and the same analysis of the vertical forces on A gives $T_{\mathrm{B} \text { on } \mathrm{A}}=\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=m_{\mathrm{A}} g$. Since $T_{\mathrm{A} \text { on } \mathrm{B}}=T_{\mathrm{B} \text { on } \mathrm{A}}$, we have $\left(f_{\mathrm{S}}\right)_{\mathrm{B}}=m_{\mathrm{A}} g$, so

$$
\mu_{\mathrm{s}} m_{\mathrm{B}} g=m_{\mathrm{A}} g \quad \Rightarrow \quad m_{\mathrm{A}}=\mu_{\mathrm{s}} m_{\mathrm{B}}=(0.60)(20 \mathrm{~kg})=12 \mathrm{~kg}
$$

7.9. Model: Model the car and the truck as particles denoted by the symbols $C$ and $T$, respectively. Denote the surface of the ground by the symbol S.

## Visualize:

## Pictorial representation



Solve: (a) The $x$-component of Newton's second law for the car gives

$$
\sum\left(F_{\text {on } \mathrm{C}}\right)_{x}=F_{\mathrm{S} \text { on } \mathrm{C}}-F_{\mathrm{T} \text { on } \mathrm{C}}=m_{\mathrm{C}} a_{\mathrm{C}}
$$

The $x$-component of Newton's second law for the truck gives

$$
\sum\left(F_{\text {on } \mathrm{T}}\right)_{x}=F_{\mathrm{C} \text { on } \mathrm{T}}=m_{\mathrm{T}} a_{\mathrm{T}}
$$

Using $a_{\mathrm{C}}=a_{\mathrm{T}} \equiv a$ and $F_{\mathrm{T} \text { on } \mathrm{C}}=F_{\mathrm{C} \text { on } \mathrm{T}}$, we get

$$
\left(F_{\mathrm{Con} \mathrm{~S}}-F_{\mathrm{ConT}}\right)\left(\frac{1}{m_{\mathrm{C}}}\right)=a \text { and }\left(F_{\mathrm{C} \text { on T }}\right)\left(\frac{1}{m_{\mathrm{T}}}\right)=a
$$

Combining these two equations,

$$
\begin{gathered}
\left(F_{\mathrm{ConS}}-F_{\mathrm{ConT}}\right)\left(\frac{1}{m_{\mathrm{C}}}\right)=\left(F_{\mathrm{ConT}}\right)\left(\frac{1}{m_{\mathrm{T}}}\right) \Rightarrow F_{\mathrm{ConT}}\left(\frac{1}{m_{\mathrm{C}}}+\frac{1}{m_{\mathrm{T}}}\right)=\left(F_{\mathrm{C} \text { on } \mathrm{S}}\right)\left(\frac{1}{m_{\mathrm{C}}}\right) \\
F_{\mathrm{ConT}=}=\left(F_{\mathrm{C} \text { on } \mathrm{S}}\right)\left(\frac{m_{\mathrm{T}}}{m_{\mathrm{C}}+m_{\mathrm{T}}}\right)=(4500 \mathrm{~N})\left(\frac{2000 \mathrm{~kg}}{1000 \mathrm{~kg}+2000 \mathrm{~kg}}\right)=3000 \mathrm{~N}
\end{gathered}
$$

(b) Due to Newton's third law, $F_{\mathrm{T} \text { on } \mathrm{C}}=3000 \mathrm{~N}$.
7.10. Model: The blocks are to be modeled as particles and denoted as 1,2 , and 3 . The surface is frictionless and along with the earth it is a part of the environment. The three blocks are our three systems of interest.

## Visualize:

## Pictorial representation



The force applied on block 1 is $F_{\text {A on } 1}=12 \mathrm{~N}$. The acceleration for all the blocks is the same and is denoted by $a$.
Solve: (a) Newton's second law for the three blocks along the $x$-direction is

$$
\sum\left(F_{\text {on } 1}\right)_{x}=F_{\mathrm{A} \text { on } 1}-F_{2 \text { on } 1}=m_{1} a, \quad \sum\left(F_{\text {on } 2}\right)_{x}=F_{1 \text { on } 2}-F_{3 \text { on } 2}=m_{2} a, \quad \sum\left(F_{\text {on } 3}\right)_{x}=F_{2 \text { on } 3}=m_{3} a
$$

Summing these three equations and using Newton's third law $\left(F_{2 \text { on } 1}=F_{1 \text { on } 2}\right.$ and $F_{3 \text { on } 2}=F_{2 \text { on 3 }}$ ), we get

$$
F_{\mathrm{A} \text { on } 1}=\left(m_{1}+m_{2}+m_{3}\right) a \Rightarrow(12 \mathrm{~N})=(1 \mathrm{~kg}+2 \mathrm{~kg}+3 \mathrm{~kg}) a \Rightarrow a=2 \mathrm{~m} / \mathrm{s}^{2}
$$

Using this value of $a$, the force equation for block 3 gives

$$
F_{2 \text { on } 3}=m_{3} a=(3 \mathrm{~kg})\left(2 \mathrm{~m} / \mathrm{s}^{2}\right)=6 \mathrm{~N}
$$

(b) Substituting into the force equation on block 1 ,

$$
12 \mathrm{~N}-F_{2 \text { on } 1}=(1 \mathrm{~kg})\left(2 \mathrm{~m} / \mathrm{s}^{2}\right) \Rightarrow F_{2 \text { on } 1}=10 \mathrm{~N}
$$

Assess: Because all three blocks are pushed forward by a force of 12 N , the value of 10 N for the force that the 2 kg block exerts on the 1 kg block seems reasonable.
7.11. Model: We treat the two objects of interest, the block (B) and steel cable (C), like particles. The motion of these objects is governed by the constant-acceleration kinematic equations. The horizontal component of the external force is 100 N .
Visualize:

## Pictorial representation



Solve: Using $v_{1 x}^{2}=v_{0 x}^{2}+2 a_{x}\left(x_{1}-x_{0}\right)$, we find

$$
(4.0 \mathrm{~m} / \mathrm{s})^{2}=0 \mathrm{~m}^{2} / \mathrm{s}^{2}+2 a_{x}(2.0 \mathrm{~m}) \Rightarrow a_{x}=4.0 \mathrm{~m} / \mathrm{s}^{2}
$$

From the free-body diagram on the block:

$$
\sum\left(F_{\text {on } \mathrm{B}}\right)_{x}=\left(F_{\mathrm{C} \text { on } \mathrm{B}}\right)_{x}=m_{\mathrm{B}} a_{x} \quad \Rightarrow \quad\left(F_{\mathrm{C} \text { on } \mathrm{B}}\right)_{x}=(20 \mathrm{~kg})\left(4.0 \mathrm{~m} / \mathrm{s}^{2}\right)=80 \mathrm{~N}
$$

Also, according to Newton's third law $\left(F_{\mathrm{B} \text { on } \mathrm{C}}\right)_{x}=\left(F_{\mathrm{C} \text { on } \mathrm{B}}\right)_{x}=80 \mathrm{~N}$. Applying Newton's second law to the cable gives

$$
\sum\left(F_{\text {on C }}\right)_{x}=\left(F_{\text {ext }}\right)_{x}-\left(F_{\mathrm{B} \text { on C }}\right)_{x}=m_{\mathrm{C}} a_{x} \Rightarrow 100 \mathrm{~N}-80 \mathrm{~N}=m_{\mathrm{C}}\left(4.0 \mathrm{~m} / \mathrm{s}^{2}\right) \Rightarrow m_{\mathrm{C}}=5.0 \mathrm{~kg}
$$

## Section 7.4 Ropes and Pulleys

7.12. Model: The man ( M ) and the block $(\mathrm{B})$ are interacting with each other through a rope. We will assume the pulley to be frictionless, which implies that the tension in the rope is the same on both sides of the pulley. The system is the man and the block.
Visualize:

## Pictorial representation



Solve: Clearly the entire system remains in equilibrium since $m_{\mathrm{B}}>m_{\mathrm{M}}$. The block would move downward but it is already on the ground. From the free-body diagrams, we can write Newton's second law in the vertical direction as

$$
\sum\left(F_{\text {on M }}\right)_{y}=T_{\mathrm{R} \text { on } \mathrm{M}}-\left(F_{\mathrm{G}}\right)_{\mathrm{M}}=0 \mathrm{~N} \Rightarrow T_{\mathrm{R} \text { on } \mathrm{M}}=\left(F_{\mathrm{G}}\right)_{\mathrm{M}}=(60 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=590 \mathrm{~N}
$$

Since the tension is the same on both sides, $T_{\mathrm{B} \text { on } \mathrm{R}}=T_{\mathrm{M} \text { on R }}=T=590 \mathrm{~N}$.
7.13. Model: The two ropes and the two blocks ( A and B ) will be treated as particles.

## Visualize:

Free-body diagrams


Solve: (a) The two blocks and two ropes form a combined system of total mass $M=2.5 \mathrm{~kg}$. This combined system is accelerating upward at $a=3.0 \mathrm{~m} / \mathrm{s}^{2}$ under the influence of a force $F$ and the gravitational force $-M g \hat{j}$. Newton's second law applied to the combined system gives

$$
\left(F_{\text {net }}\right)_{y}=F-M g=M a \Rightarrow F=M(a+g)=(2.5 \mathrm{~kg})\left(3.0 \mathrm{~m} / \mathrm{s}^{2}+9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=32 \mathrm{~N}
$$

(b) The ropes are not massless. We must consider both the blocks and the ropes as systems. The force $F$ acts only on block A because it does not contact the other objects. We can proceed to apply the $y$-component of Newton's second law to each system, starting at the top. Each object accelerates upward at $a=3.0 \mathrm{~m} / \mathrm{s}^{2}$. For block A,

$$
\left(F_{\text {net on } \mathrm{A}}\right)_{y}=F-m_{\mathrm{A}} g-T_{1 \text { on } \mathrm{A}}=m_{\mathrm{A}} a \Rightarrow T_{1 \text { on } \mathrm{A}}=F-m_{\mathrm{A}}(a+g)=19 \mathrm{~N}
$$

(c) Applying Newton's second law to rope 1 gives

$$
\left(F_{\text {net on } 1}\right)_{y}=T_{\mathrm{A} \text { on } 1}-m_{1} g-T_{\mathrm{B} \text { on } 1}=m_{1} a
$$

where $\vec{T}_{\mathrm{A} \text { on } 1}$ and $\vec{T}_{1 \text { on A }}$ are an action/reaction pair. But, because the rope has mass, the two tension forces $\vec{T}_{\mathrm{A}}$ on 1 and $\vec{T}_{\mathrm{B} \text { on } 1}$ are not the same. The tension at the lower end of rope 1 , where it connects to B , is

$$
T_{\mathrm{B} \text { on } 1}=T_{\mathrm{A} \text { on } 1}-m_{1}(a+g)=16 \mathrm{~N}
$$

(d) We can continue to repeat this procedure, noting from Newton's third law that

$$
T_{1 \text { on } \mathrm{B}}=T_{\mathrm{B} \text { on } 1} \text { and } T_{2 \text { on } \mathrm{B}}=T_{\mathrm{B} \text { on } 2}
$$

Newton's second law applied to block B is

$$
\left(F_{\text {net on } \mathrm{B}}\right)_{y}=T_{1 \text { on } \mathrm{B}}-m_{\mathrm{B}} g-T_{2 \text { on } \mathrm{B}}=m_{\mathrm{B}} a \Rightarrow T_{2 \text { on } \mathrm{B}}=T_{1 \text { on } \mathrm{B}}-m_{\mathrm{B}}(a+g)=3.2 \mathrm{~N}
$$

7.14. Model: Together the carp (C) and the trout $(\mathrm{T})$ make up the system that will be represented through the particle model. The fishing rod line (R) is assumed to be massless.

## Visualize:

## Pictorial representation



Known
$m_{\mathrm{C}}=1.5 \mathrm{~kg}$
$m_{\mathrm{T}}=3.0 \mathrm{~kg}$
$T_{2}=60 \mathrm{~N}$

$$
\frac{\text { Find }}{T_{1}\left(F_{\Gamma_{\mathrm{F}}}\right)_{\mathrm{T}}\left(F_{\Gamma_{\mathrm{F}}}\right)_{\Gamma} T_{7}}
$$



Solve: Jimmy's pull $T_{2}=60 \mathrm{~N}$ is larger than the total weight of the fish, so they accelerate upward. They are tied together, so each fish has the same acceleration $a$. Applying Newton's second law along the $y$-direction for the carp and the trout gives

$$
\sum\left(F_{\mathrm{on} \mathrm{C}}\right)_{y}=T_{2}-T_{1}-\left(F_{\mathrm{G}}\right)_{\mathrm{C}}=m_{\mathrm{C}} a \Rightarrow \sum\left(F_{\mathrm{on} \mathrm{~T}}\right)_{y}=T_{1}-\left(F_{\mathrm{G}}\right)_{\mathrm{T}}=m_{\mathrm{T}} a
$$

Adding these two equations gives

$$
a=\frac{T_{2}-\left(F_{\mathrm{G}}\right)_{\mathrm{C}}-\left(F_{\mathrm{G}}\right)_{\mathrm{T}}}{\left(m_{\mathrm{C}}+m_{\mathrm{T}}\right)}=\frac{60 \mathrm{~N}-(1.5 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)-(3.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{1.5 \mathrm{~kg}+3.0 \mathrm{~kg}}=3.533 \mathrm{~m} / \mathrm{s}^{2}
$$

Substituting this value of acceleration back into the force equation for the trout, we find that

$$
T_{1}=m_{\mathrm{T}}(a+g)=(3.0 \mathrm{~kg})\left(3.533 \mathrm{~m} / \mathrm{s}^{2}+9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=40 \mathrm{~N}
$$

$$
\left(F_{\mathrm{G}}\right)_{\mathrm{T}}=m_{\mathrm{T}} g=(3.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=29 \mathrm{~N} \quad \Rightarrow \quad\left(F_{\mathrm{G}}\right)_{\mathrm{C}}=m_{\mathrm{C}} g=(1.5 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=15 \mathrm{~N}
$$

Thus, $T_{2}>T_{1}>\left(F_{\mathrm{G}}\right)_{\mathrm{T}}>\left(F_{\mathrm{G}}\right)_{\mathrm{C}}$.
7.15. Model: The block of ice (I) is a particle and so is the rope (R) because it is not massless. We must therefore consider both the block of ice and the rope as objects in the system.

## Visualize:



Solve: (a) The force $\vec{F}_{\text {ext }}$ acts only on the rope. Since the rope and the ice block move together, they have the same acceleration. Also because the rope has mass, $F_{\text {ext }}$ on the front end of the rope is not the same as $F_{\text {Ion R }}$ that acts on the rear end of the rope. Applying Newton's second law along the $x$-axis to the ice block and the rope gives

$$
\sum\left(F_{\text {on I }}\right)_{x}=\left(F_{\mathrm{R} \text { on I }}\right)_{x}=m_{\mathrm{I}} a=(10 \mathrm{~kg})\left(2.0 \mathrm{~m} / \mathrm{s}^{2}\right)=20 \mathrm{~N}
$$

(b) Applying Newton's second law to the rope gives

$$
\sum\left(F_{\text {on R }}\right)_{x}=\left(F_{\mathrm{ext}}\right)_{x}-\left(F_{\mathrm{I} \text { on R }}\right)_{x}=m_{\mathrm{R}} a \quad \Rightarrow \quad\left(F_{\mathrm{ext}}\right)_{x}=\left(F_{\mathrm{R} \text { on I }}\right)_{x}+m_{\mathrm{R}} a=20 \mathrm{~N}+(0.500 \mathrm{~kg})\left(2.0 \mathrm{~m} / \mathrm{s}^{2}\right)=21 \mathrm{~N}
$$

7.16. Model: The hanging block and the rail car are objects in the systems.

## Visualize:

Pictorial representation



Solve: The mass of the rope is very small in comparison to the $2000-\mathrm{kg}$ block, so we assume a massless rope. In this case, the forces $\vec{T}_{1}$ and $\vec{T}_{1}^{\prime}$ act as if they are an action/reaction pair. The hanging block is in static equilibrium, with $\vec{F}_{\text {net }}=0 \mathrm{~N}$, so $T_{1}^{\prime}=m_{\text {block }} g=(2000 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=19,600 \mathrm{~N}$. The rail car with the pulley is also in static equilibrium, so $T_{2}+T_{3}-T_{1}=0 \mathrm{~N}$. Notice how the tension force in the cable pulls both the top and bottom of the pulley to the right. Now, $T_{1}=T_{1}^{\prime}=19,600 \mathrm{~N}$ by Newton's third law. Also, the cable tension is $T_{2}=T_{3}=T$. Thus, $T=\frac{1}{2} T_{1}^{\prime}=9800 \mathrm{~N}$.

### 7.17. Visualize:



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Solve: The rope is treated as two $1.0-\mathrm{kg}$ interacting objects. At the midpoint of the rope, the rope has a tension $T_{\mathrm{B} \text { on } \mathrm{T}}=T_{\mathrm{T} \text { on } \mathrm{B}} \equiv T$. Apply Newton's first law to the bottom half of the rope to find $T$.

$$
\left(F_{\text {net }}\right)_{y}=0=T-\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \Rightarrow T=m_{\mathrm{B}} g=(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=9.8 \mathrm{~N}
$$

Assess: 9.8 N is half the gravitational force on the whole rope. This is reasonable since the top half is holding up the bottom half of the rope against gravity.
7.18. Model: The cat and dog are modeled as two blocks in the pictorial representation below. The rope is assumed to be massless. The two points (knots) where the blocks are attached to the rope and the two hanging blocks form a system. These four objects are treated at particles, form the system, and are in static equilibrium.

## Visualize:

## Pictorial representation




Solve: (a) We consider both the two hanging blocks and the two knots. The blocks are in static equilibrium with $\vec{F}_{\text {net }}=0 \mathrm{~N}$. Note that there are three action/reaction pairs. For Block 1 and Block 2, $\vec{F}_{\text {net }}=0 \mathrm{~N}$ and we have

$$
T_{4}^{\prime}=\left(F_{\mathrm{G}}\right)_{1}=m_{1} g, \quad T_{5}^{\prime}=\left(F_{\mathrm{G}}\right)_{2}=m_{2} g
$$

By Newton's third law:

$$
T_{4}=T_{4}^{\prime}=m_{1} g, \quad T_{5}=T_{5}^{\prime}=m_{2} g
$$

The knots are also in equilibrium. Newton's law applied to the left knot is

$$
\left(F_{\text {net }}\right)_{x}=T_{2}-T_{1} \cos \theta_{1}=0 \mathrm{~N} \Rightarrow\left(F_{\text {net }}\right)_{y}=T_{1} \sin \theta_{1}-T_{4}=T_{1} \sin \theta_{1}-m_{1} g=0 \mathrm{~N}
$$

The $y$-equation gives $T_{1}=m_{1} g / \sin \theta_{1}$. Substitute this into the $x$-equation to find

$$
T_{2}=\frac{m_{1} g \cos \theta_{1}}{\sin \theta_{1}}=\frac{m_{1} g}{\tan \theta_{1}}
$$

Newton's law applied to the right knot is

$$
\left(F_{\text {net }}\right)_{x}=T_{3} \cos \theta_{3}-T_{2}^{\prime}=0 \mathrm{~N} \Rightarrow\left(F_{\text {net }}\right)_{y}=T_{3} \sin \theta_{3}-T_{5}=T_{3} \sin \theta_{3}-m_{2} g=0 \mathrm{~N}
$$

These can be combined just like the equations for the left knot to give

$$
T_{2}^{\prime}=\frac{m_{2} g \cos \theta_{3}}{\sin \theta_{3}}=\frac{m_{2} g}{\tan \theta_{3}}
$$

But the forces $\vec{T}_{2}$ and $\vec{T}_{2}^{\prime}$ are an action/reaction pair, so $T_{2}=T_{2}^{\prime}$. Therefore,

$$
\frac{m_{1} g}{\tan \theta_{1}}=\frac{m_{2} g}{\tan \theta_{3}} \Rightarrow \tan \theta_{3}=\frac{m_{2}}{m_{1}} \tan \theta_{1} \Rightarrow \theta_{3}=\tan ^{-1}\left[2 \tan \left(20^{\circ}\right)\right]=36^{\circ}
$$

We can now use the $y$-equation for the right knot to find $T_{3}=m_{2} g / \sin \theta_{3}=67 \mathrm{~N}$.
7.19. (a) Visualize: The upper magnet is labeled $U$ and the lower magnet $L$. Each magnet exerts a long-range magnetic force on the other. Each magnet and the table exert a contact force (normal force) on each other. In addition, the table experiences a normal force due to the surface.

## Interaction diagram



$$
\begin{array}{ll}
\mathrm{U}=\text { Upper magnet } & \text { Known } \\
=\text { Lower magnet } & \left(F_{\mathrm{G}}\right)_{\mathrm{T}}=20.0 \mathrm{~N} \\
\mathrm{~T}=\text { Table } & \left(F_{\mathrm{G}}\right)_{\mathrm{U}}=2.0 \mathrm{~N} \\
\mathrm{~S}=\text { Surface } & \left(F_{\mathrm{G}}\right)_{\mathrm{L}}=2.0 \mathrm{~N} \\
\mathrm{EE}=\text { Entire Earth } & F_{\mathrm{U} \text { on } \mathrm{L}}=3\left(F_{\mathrm{G}}\right)_{\mathrm{L}}
\end{array}
$$


(b) Solve: Each object is in static equilibrium with $F_{\text {net }}=0$. Start with the lower magnet. Because $F_{\mathrm{U} \text { on } \mathrm{L}}=3\left(F_{\mathrm{G}}\right)_{\mathrm{L}}=6.0 \mathrm{~N}$, equilibrium requires $n_{\mathrm{T} \text { on } \mathrm{L}}=4.0 \mathrm{~N}$. For the upper magnet, $F_{\mathrm{L} \text { on } \mathrm{U}}=F_{\mathrm{U} \text { on } \mathrm{L}}=6.0 \mathrm{~N}$ because these are an action/reaction pair. Equilibrium for the upper magnet requires $n_{T \text { on } U}=8.0 \mathrm{~N}$. For the table, the action/reaction pairs are $n_{\mathrm{L} \text { on } \mathrm{T}}=n_{\mathrm{T} \text { on } \mathrm{L}}=4.0 \mathrm{~N}$ and $n_{\mathrm{U} \text { on } \mathrm{T}}=n_{\mathrm{T} \text { on } \mathrm{U}}=8.0 \mathrm{~N}$. The table's gravitational force is $\left(F_{\mathrm{G}}\right)_{\mathrm{T}}=20 \mathrm{~N}$, so $n_{\mathrm{S} \text { on } \mathrm{T}}=24 \mathrm{~N}$ for the table to be in equilibrium. Summarizing, we have

| Upper magnet | Table | Lower magnet |
| :---: | :---: | :---: |
| $\left(F_{\mathrm{G}}\right)_{\mathrm{U}}=2.0 \mathrm{~N}$ | $\left(F_{\mathrm{G}}\right)_{\mathrm{T}}=20 \mathrm{~N}$ | $\left(F_{\mathrm{G}}\right)_{\mathrm{L}}=2.0 \mathrm{~N}$ |
| $n_{\mathrm{T} \text { on } \mathrm{U}}=8.0 \mathrm{~N}$ | $n_{\mathrm{U} \text { on } \mathrm{T}}=8.0 \mathrm{~N}$ | $n_{\mathrm{T} \text { on } \mathrm{L}}=4.0 \mathrm{~N}$ |
| $F_{\mathrm{L} \text { on } \mathrm{U}}=6.0 \mathrm{~N}$ | $n_{\mathrm{L} \text { on } \mathrm{T}}=4.0 \mathrm{~N}$ | $F_{\mathrm{U} \text { on } \mathrm{L}}=6.0 \mathrm{~N}$ |
|  | $n_{\mathrm{S} \text { on } \mathrm{T}}=24 \mathrm{~N}$ |  |

Assess: The result $n_{\text {S on } T}=24 \mathrm{~N}$ makes sense. The combined gravitational force on the table and two magnets is 24 N . Because the table is in equilibrium, the upward normal force of the surface has to exactly balance the total gravitational force on the table and magnets.
7.20. Model: The astronaut and the satellite, the two objects in our system, will be treated as particles. Visualize:

## Pictorial representation



> | Known |
| :--- |
| $m_{\mathrm{A}}=80 \mathrm{~kg}$ |
| $x_{0 \mathrm{~A}}=x_{0 \mathrm{~S}}=0 \quad m_{\mathrm{S}}=640 \mathrm{~kg}$ |
| $v_{0 \mathrm{~A}}=0$ |
| $F_{\mathrm{AS}}=0$ |
| $t_{1}=0.50 \mathrm{~s}$ S |
| $\mathrm{S}_{\mathrm{S} \text { on } \mathrm{A}}=100 \mathrm{~N}$ |
| $t_{2}=60.0 \mathrm{~s}$ |

$\frac{\text { Find }}{x_{2 \mathrm{~A}}-x_{2 S}}$

Solve: The astronaut and the satellite accelerate in opposite directions for 0.50 s . The force on the satellite and the force on the astronaut are an action/reaction pair, so both have a magnitude of 100 N . Newton's second law for the satellite along the $x$-direction gives

$$
\sum\left(F_{\mathrm{on} \mathrm{~S}}\right)_{x}=F_{\mathrm{A} \text { on } \mathrm{S}}=m_{\mathrm{S}} a_{\mathrm{S}} \Rightarrow a_{\mathrm{S}}=\frac{F_{\mathrm{A} \text { on } \mathrm{S}}}{m_{\mathrm{S}}}=\frac{-(100 \mathrm{~N})}{640 \mathrm{~kg}}=-0.156 \mathrm{~m} / \mathrm{s}^{2}
$$

Newton's second law for the astronaut along the $x$-direction is

$$
\sum\left(F_{\text {on } \mathrm{A}}\right)_{x}=F_{\mathrm{S} \text { on } \mathrm{A}}=m_{\mathrm{A}} a_{\mathrm{A}} \Rightarrow a_{\mathrm{A}}=\frac{F_{\mathrm{S} \text { on } \mathrm{A}}}{m_{\mathrm{A}}}=\frac{F_{\mathrm{A} \text { on } \mathrm{S}}}{m_{\mathrm{A}}}=\frac{100 \mathrm{~N}}{80 \mathrm{~kg}}=1.25 \mathrm{~m} / \mathrm{s}^{2}
$$

Let us first calculate the positions and velocities of the astronaut and the satellite at $t_{1}=0.50 \mathrm{~s}$ under the accelerations $a_{\mathrm{A}}$ and $a_{\mathrm{S}}$ :

$$
\begin{gathered}
x_{1 \mathrm{~A}}=x_{0 \mathrm{~A}}+v_{0 \mathrm{~A}}\left(t_{1}-t_{0}\right)+\frac{1}{2} a_{\mathrm{A}}\left(t_{1}-t_{0}\right)^{2}=0 \mathrm{~m}+0 \mathrm{~m}+\frac{1}{2}\left(1.25 \mathrm{~m} / \mathrm{s}^{2}\right)(0.50 \mathrm{~s}-0.00 \mathrm{~s})^{2}=0.156 \mathrm{~m} \\
x_{1 \mathrm{~S}}=x_{0 \mathrm{~S}}+v_{0 \mathrm{~S}}\left(t_{1}-t_{0}\right)+\frac{1}{2} a_{\mathrm{S}}\left(t_{1}-t_{0}\right)^{2}=0 \mathrm{~m}+0 \mathrm{~m}+\frac{1}{2}\left(-0.156 \mathrm{~m} / \mathrm{s}^{2}\right)(0.50 \mathrm{~s}-0.00 \mathrm{~s})^{2}=-0.020 \mathrm{~m} \\
v_{1 \mathrm{~A}}=v_{0 \mathrm{~A}}+a_{\mathrm{A}}\left(t_{1}-t_{0}\right)=0 \mathrm{~m} / \mathrm{s}+\left(1.25 \mathrm{~m} / \mathrm{s}^{2}\right)(0.50 \mathrm{~s}-0.00 \mathrm{~s})=0.625 \mathrm{~m} / \mathrm{s} \\
v_{1 \mathrm{~S}}=v_{0 \mathrm{~S}}+a_{\mathrm{S}}\left(t_{1}-t_{0}\right)=0 \mathrm{~m} / \mathrm{s}+\left(-0.156 \mathrm{~m} / \mathrm{s}^{2}\right)(0.5 \mathrm{~s}-0.00 \mathrm{~s})=-0.078 \mathrm{~m} / \mathrm{s}
\end{gathered}
$$

With $x_{1 \mathrm{~A}}$ and $x_{1 \mathrm{~S}}$ as initial positions, $v_{1 \mathrm{~A}}$ and $v_{1 \mathrm{~S}}$ as initial velocities, and zero accelerations, we can now obtain the new positions at $\left(t_{2}-t_{1}\right)=59.5 \mathrm{~s}$ :

$$
\begin{aligned}
& x_{2 \mathrm{~A}}=x_{1 \mathrm{~A}}+v_{1 \mathrm{~A}}\left(t_{2}-t_{1}\right)=0.156 \mathrm{~m}+(0.625 \mathrm{~m} / \mathrm{s})(59.5 \mathrm{~s})=37.34 \mathrm{~m} \\
& x_{2 \mathrm{~S}}=x_{1 \mathrm{~S}}+v_{1 \mathrm{~S}}\left(t_{2}-t_{1}\right)=-0.02 \mathrm{~m}+(-0.078 \mathrm{~m} / \mathrm{s})(59.5 \mathrm{~s})=-4.66 \mathrm{~m}
\end{aligned}
$$

Thus the astronaut and the satellite are $x_{2 \mathrm{~A}}-x_{2 \mathrm{~S}}=(37.34 \mathrm{~m})-(-4.66 \mathrm{~m})=42 \mathrm{~m}$ apart.
7.21. Model: The block (B) and the steel cable (C), the two objects in the system, are modeled as particles and their motion is determined by the constant-acceleration kinematic equations.

## Visualize:

Pictorial representation


$$
\begin{aligned}
& \text { Known } \\
& \hline m_{\mathrm{B}}=20 \mathrm{~kg} \\
& \left(\vec{F}_{\text {ext }}\right)_{x}=100 \mathrm{~N} \\
& x_{0}=v_{0 x}=t_{0}=0 \\
& v_{1 x}=4.0 \mathrm{~m} / \mathrm{s} \\
& t_{1}=2.0 \mathrm{~s}
\end{aligned}
$$

$$
\frac{\text { Find }}{\left(\vec{F}_{\text {ext }}\right)_{x}-\left(\vec{F}_{\mathrm{B} \text { on } \mathrm{C}}\right)}
$$


Block

Cable

Solve: Using $v_{1 x}=v_{0 x}+a_{x}\left(t_{1}-t_{0}\right)$,

$$
4.0 \mathrm{~m} / \mathrm{s}=0 \mathrm{~m} / \mathrm{s}+a_{x}(2.0 \mathrm{~s}-0.0 \mathrm{~s}) \Rightarrow a_{x}=2.0 \mathrm{~m} / \mathrm{s}^{2}
$$

Newton's second law along the $x$-direction for the block gives

$$
\sum\left(F_{\text {on B }}\right)_{x}=\left(F_{\mathrm{C} \text { on B }}\right)_{x}=m_{\mathrm{B}} a_{x}=(20 \mathrm{~kg})\left(2.0 \mathrm{~m} / \mathrm{s}^{2}\right)=40 \mathrm{~N}
$$

$\left(F_{\text {ext }}\right)_{x}$ acts on the right end of the cable and $\left(F_{\mathrm{B} \text { on } \mathrm{C}}\right)_{x}$ acts on the left end. According to Newton's third law, $\left(F_{\mathrm{B} \text { on } \mathrm{C}}\right)_{x}=\left(F_{\mathrm{C} \text { on B }}\right)_{x}=40 \mathrm{~N}$. The difference in the horizontal component of the tension between the two ends of the cable is thus

$$
\left(F_{\mathrm{ext}}\right)_{x}-\left(F_{\mathrm{B} \text { on } \mathrm{C}}\right)_{x}=100 \mathrm{~N}-40 \mathrm{~N}=60 \mathrm{~N}
$$

7.22. Model: The gliders and spring form the system and are modeled as particles. Because the spring is compressed, it may be modeled as a rigid rod, so the three objects are constrained to have the same acceleration.

## Visualize:

## Pictorial representation


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Solve: By Newton's third law, we know the action/reaction forces are of equal magnitude but point in the opposite direction, so they cancel when considering the entire system. Using Newton's second law in the $x$-direction, the acceleration of the system is

$$
\sum(F)_{x}=F_{\mathrm{ext}}=\left(m_{\mathrm{A}}+m_{\mathrm{Sp}}+m_{\mathrm{B}}\right) a \Rightarrow a=\frac{F_{\mathrm{ext}}}{m_{\mathrm{A}}+m_{\mathrm{Sp}}+m_{\mathrm{B}}}=\frac{6.0 \mathrm{~N}}{0.40 \mathrm{~kg}+0.20 \mathrm{~kg}+0.60 \mathrm{~kg}}=5.0 \mathrm{~m} / \mathrm{s}^{2}
$$

Applying Newton's second law to glider A gives

$$
\sum(F)_{x}=F_{\mathrm{ext}}-\left(\vec{F}_{\mathrm{Sp} \mathrm{on} \mathrm{~A}}\right)_{x}=m_{\mathrm{A}} a \Rightarrow\left(\vec{F}_{\mathrm{Sp} \mathrm{on} \mathrm{~A}}\right)_{x}=F_{\mathrm{ext}}-m_{\mathrm{A}} a=6.0 \mathrm{~N}-(0.40 \mathrm{~kg})\left(5.0 \mathrm{~m} / \mathrm{s}^{2}\right)=4.0 \mathrm{~N}
$$

Applying Newton's second law to glider B gives

$$
\sum(F)_{x}=\left(\vec{F}_{\text {Sp on B }}\right)_{x}=m_{\mathrm{B}} a=(0.60 \mathrm{~kg})\left(5.0 \mathrm{~m} / \mathrm{s}^{2}\right)=3.0 \mathrm{~N}
$$

Applying Newton's second law to the spring in the $y$-direction, and using the fact that $\left(\vec{F}_{\mathrm{Sp} \mathrm{on} \mathrm{A}}\right)_{y}=\left(\vec{F}_{\mathrm{Sp} \text { on B }}\right)_{y}$ by symmetry, we find

$$
\begin{aligned}
\sum(F)_{y} & =\left(\vec{F}_{\text {Sp on A }}\right)_{y}+\left(\vec{F}_{\mathrm{Sp} \mathrm{on} \mathrm{~B}}\right)_{y}-\left(F_{\mathrm{G}}\right)_{\mathrm{Sp}}=0 \Rightarrow 2\left(\vec{F}_{\mathrm{Sp} \mathrm{on} \mathrm{~B}}\right)_{y}=m_{\mathrm{Sp}}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \\
\left(\vec{F}_{\text {Sp on A }}\right)_{y} & =\left(\vec{F}_{\mathrm{Sp} \text { on } \mathrm{B}}\right)_{y}=\frac{1}{2}(0.20 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=0.98 \mathrm{~N}
\end{aligned}
$$

Adding the $x$ - and $y$-components in quadrature gives the total force exerted by the spring on each block:

$$
\begin{array}{ll}
\text { glider } \mathrm{A}: & F_{\mathrm{Sp} \text { on } \mathrm{A}}=\sqrt{(4.0 \mathrm{~N})^{2}+(0.98 \mathrm{~N})^{2}}=4.1 \mathrm{~N} \\
\text { glider } \mathrm{B}: & F_{\mathrm{Sp} \text { on } \mathrm{B}}=\sqrt{(3.0 \mathrm{~N})^{2}+(0.98 \mathrm{~N})^{2}}=3.2 \mathrm{~N}
\end{array}
$$

Assess: The result seems reasonable because more force is exerted on glider A by the spring than on glider B, as expected. The force exerted on glider A by the spring is, by Newton's third law, the force that must accelerate the spring + glider B, whereas the force exerted by the spring on glider B only has to accelerate glider B.
7.23. Model: Sled A, sled B, and the dog (D) are treated like particles in the model. The rope is considered to be massless.

## Visualize:

## Pictorial representation



$$
\frac{\text { Find }}{T_{2}}
$$



Solve: The acceleration constraint is $\left(a_{\mathrm{A}}\right)_{x}=\left(a_{\mathrm{B}}\right)_{x}=a_{x}$. Newton's second law applied to sled A gives

$$
\begin{aligned}
& \sum\left(\vec{F}_{\text {on A }}\right)_{y}=n_{\mathrm{A}}-\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{A}}=\left(F_{\mathrm{G}}\right)_{\mathrm{A}}=m_{\mathrm{A}} g \\
& \sum\left(\vec{F}_{\text {on A }}\right)_{x}=T_{1 \text { on } \mathrm{A}}-f_{\mathrm{A}}=m_{\mathrm{A}} a_{x}
\end{aligned}
$$

Using $f_{\mathrm{A}}=\mu_{\mathrm{k}} n_{\mathrm{A}}$, the $x$-equation yields

$$
T_{1 \text { on A }}-\mu_{\mathrm{k}} n_{\mathrm{A}}=m_{\mathrm{A}} a_{x} \Rightarrow 150 \mathrm{~N}-(0.10)(100 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=(100 \mathrm{~kg}) a_{x} \Rightarrow a_{x}=0.52 \mathrm{~m} / \mathrm{s}^{2}
$$

Newton's second law applied to sled B gives

$$
\begin{aligned}
& \sum\left(\vec{F}_{\text {on B }}\right)_{y}=n_{\mathrm{B}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{B}}=\left(F_{\mathrm{G}}\right)_{\mathrm{B}}=m_{\mathrm{B}} g \\
& \sum\left(\vec{F}_{\text {on B }}\right)_{x}=T_{2}-T_{1 \text { on } \mathrm{B}}-f_{\mathrm{B}}=m_{\mathrm{B}} a_{x}
\end{aligned}
$$

$T_{1 \text { on B }}$ and $T_{1 \text { on A }}$ act as if they are an action/reaction pair, so $T_{1 \text { on } \mathrm{B}}=150 \mathrm{~N}$. Using $f_{\mathrm{B}}=\mu_{\mathrm{k}} n_{\mathrm{B}}=(0.10)(80 \mathrm{~kg})$ $\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=78.4 \mathrm{~N}$, we find

$$
T_{2}-150 \mathrm{~N}-78.4 \mathrm{~N}=(80 \mathrm{~kg})\left(0.52 \mathrm{~m} / \mathrm{s}^{2}\right) \Rightarrow T_{2}=270 \mathrm{~N}
$$

Thus the tension $T_{2}=2.7 \times 10^{2} \mathrm{~N}$.
7.24. Model: Consider an element of the rope $d m=\rho d y$, where $\rho=m / L$ is the mass density of the rope. Model this element as a particle.

## Visualize:

## Pictorial representation



Solve: The rope is stationary, so Newton's second law applied to the particle gives

$$
\sum\left(F_{d m}\right)_{y}=T(y)-F_{\mathrm{G}}(y)=0 \Rightarrow T(y)=F_{\mathrm{G}}(y)
$$

The force $F_{\mathrm{G}}(y)$ due to gravity is the weight of the rope below the point $y$, which is

$$
F_{\mathrm{G}}(y)=y \rho g=y(\mathrm{~m} / \mathrm{L}) g
$$

Inserting this into the expression above gives the tension: $T(y)=y m g / L$.
Assess: The result seems reasonable because $T(y)=0$ at the bottom of the rope $(y=0)$ and $T(y)=m g$ at the top of the rope $(y=L)$.
7.25. Model: The coffee mug (M) is the only object in the system, and it will be treated as a particle. The model of friction and the constant-acceleration kinematic equations will also be used.

## Visualize:

## Pictorial representation



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Solve: The mug and the car have the same velocity. If the mug does not slip, their accelerations will also be the same. Using $v_{1 x}^{2}=v_{0 x}^{2}+2 a_{x}\left(x_{1}-x_{0}\right)$, we get

$$
0 \mathrm{~m}^{2} / \mathrm{s}^{2}=(20 \mathrm{~m} / \mathrm{s})^{2}+2 a_{x}(50 \mathrm{~m}) \Rightarrow a_{x}=-4.0 \mathrm{~m} / \mathrm{s}^{2}
$$

The static force needed to stop the mug is

$$
\left(F_{\mathrm{net}}\right)_{x}=-f_{\mathrm{s}}=m a_{x}=(0.5 \mathrm{~kg})\left(-4.0 \mathrm{~m} / \mathrm{s}^{2}\right)=-2.0 \mathrm{~N} \Rightarrow f_{\mathrm{s}}=2.0 \mathrm{~N}
$$

The maximum force of static friction is

$$
\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n=\mu_{\mathrm{s}} F_{\mathrm{G}}=\mu_{\mathrm{s}} m g=(0.50)(0.50 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=2.5 \mathrm{~N}
$$

Since $\left(f_{\mathrm{s}}\right)_{\max }<\left(f_{\mathrm{s}}\right)_{\max }$, the mug does not slide.
7.26. Model: For car tires on dry concrete, the coefficient $\mu_{\mathrm{s}}$ of static friction is typically about 0.80 (see Table 6.1).

Visualize: The car and the ground are denoted by C and S , respectively.

## Pictorial representation



Solve: The car presses down against the ground with both the drive wheels (assumed to be the front wheels F , although this is not critical) and the nondrive wheels. For this car, two-thirds of the gravitational force rests on the front wheels. Physically, force $\vec{F}_{\mathrm{S} \text { on } \mathrm{C}}$ is a static friction force, so its maximum value is $\left(\vec{F}_{\mathrm{S} \text { on C }}\right)_{\max }=\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{S}} n$. The maximum acceleration of the car on the ground (or concrete surface) occurs when the static friction reaches this maximum possible value:

$$
\left(F_{\mathrm{S} \text { on C }}\right)_{\max }=\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n_{\mathrm{F}}=\mu_{\mathrm{s}}\left(F_{\mathrm{G}}\right)_{\mathrm{F}}=\mu_{\mathrm{s}}\left(\frac{2}{3} m g\right)=(0.80)\left(\frac{2}{3}\right)(1500 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=7840 \mathrm{~N}
$$

Use this force in Newton's second law to find the acceleration:

$$
a_{\max }=\frac{\left(F_{\mathrm{SonC}}\right)_{\max }}{m}=\frac{7840 \mathrm{~N}}{1500 \mathrm{~kg}}=5.2 \mathrm{~m} / \mathrm{s}^{2}
$$

7.27. Model: The starship and the shuttlecraft will be denoted as $M$ and $m$, respectively, and both will be treated as particles. We will also use the constant-acceleration kinematic equations.

## Visualize:

## Pictorial representation



Solve: (a) The tractor beam is some kind of long-range force $\vec{F}_{\mathrm{M} \mathrm{on} \mathrm{m}}$. Regardless of what kind of force it is, by Newton's third law there must be a reaction force $\vec{F}_{\mathrm{m} \text { on } \mathrm{M}}$ on the starship. As a result, both the shuttlecraft and the starship move toward each other (rather than the starship remaining at rest as it pulls the shuttlecraft in). However, the very different masses of the two crafts means that the distances they each move will also be very different. The pictorial representation shows that they meet at time $t_{1}$ when $x_{\mathrm{M} 1}=x_{\mathrm{m} 1}$. There's only one force on each craft, so Newton's second law is very simple. Furthermore, because the forces are an action/reaction pair,

$$
F_{\mathrm{M} \mathrm{on} \mathrm{~m}}=F_{\mathrm{m} \text { on M }}=F_{\text {tractor beam }}=4.0 \times 10^{4} \mathrm{~N}
$$

The accelerations of the two craft are

$$
a_{M}=\frac{F_{\mathrm{m} \text { on } \mathrm{M}}}{M}=\frac{4.0 \times 10^{4} \mathrm{~N}}{2.0 \times 10^{6} \mathrm{~kg}}=0.020 \mathrm{~m} / \mathrm{s}^{2} \text { and } a_{m}=\frac{\vec{F}_{\mathrm{M} \text { on } \mathrm{m}}}{m}=\frac{-4.0 \times 10^{4} \mathrm{~N}}{2.0 \times 10^{4} \mathrm{~kg}}=-2.0 \mathrm{~m} / \mathrm{s}^{2}
$$

Acceleration $a_{\mathrm{m}}$ is negative because the force and acceleration vectors point in the negative $x$-direction. Now we have a constant-acceleration problem in kinematics. At a later time $t_{1}$ the positions of the crafts are

$$
\begin{aligned}
& x_{\mathrm{M} 1}=x_{\mathrm{M} 0}+v_{\mathrm{M} 0}\left(t_{1}-t_{0}\right)+\frac{1}{2} a_{\mathrm{M}}\left(t_{1}-t_{0}\right)^{2}=\frac{1}{2} a_{\mathrm{M}} t_{1}^{2} \\
& x_{\mathrm{m} 1}=x_{\mathrm{m} 0}+v_{\mathrm{m} 0}\left(t_{1}-t_{0}\right)+\frac{1}{2} a_{\mathrm{m}}\left(t_{1}-t_{0}\right)^{2}=x_{\mathrm{m} 0}+\frac{1}{2} a_{\mathrm{m}} t_{1}^{2}
\end{aligned}
$$

The craft meet when $x_{\mathrm{M} 1}=x_{\mathrm{m} 1}$, so

$$
\frac{1}{2} a_{\mathrm{M}} t_{1}^{2}=x_{\mathrm{m} 0}+\frac{1}{2} a_{\mathrm{m}} t_{1}^{2} \Rightarrow t_{1}=\sqrt{\frac{2 x_{\mathrm{m} 0}}{a_{\mathrm{M}}-a_{\mathrm{m}}}}=\sqrt{\frac{2 x_{\mathrm{m} 0}}{a_{\mathrm{M}}+\left|a_{\mathrm{m}}\right|}}=\sqrt{\frac{2(10,000 \mathrm{~m})}{2.02 \mathrm{~m} / \mathrm{s}^{2}}}=99.5 \mathrm{~s}
$$

Knowing $t_{1}$, we can now find the starship's position as it meets the shuttlecraft:

$$
x_{\mathrm{M} 1}=\frac{1}{2} a_{\mathrm{M}} t_{1}^{2}=99 \mathrm{~m}
$$

The starship moves 99 m as it pulls in the shuttlecraft from 10 km away.
7.28. Model: We shall only consider horizontal forces. The head and the baseball are the two objects in our system and are treated as particles. We will also use the constant-acceleration kinematic equations.

## Visualize:



Solve: (a) The ball experiences an average acceleration of

$$
a_{B}=\frac{v_{\mathrm{B} 1}-v_{\mathrm{B} 0}}{t}=\frac{-30 \mathrm{~m} / \mathrm{s}}{1.5 \times 10^{-3} \mathrm{~s}}=-20,000 \mathrm{~m} / \mathrm{s}^{2}
$$

Insert this into Newton's second law to find the force on the baseball:

$$
F_{\mathrm{Hon} \mathrm{~B}}=m_{\mathrm{B}} a_{\mathrm{B}}=(0.14 \mathrm{~kg})\left|-20,000 \mathrm{~m} / \mathrm{s}^{2}\right|=2800 \mathrm{~N}
$$

(b) By Newton's third law, the magnitude of the force exerted by the ball on the head is the same as that exerted by the head on the ball. Thus, $F_{\mathrm{B} \text { on } \mathrm{H}}=2800 \mathrm{~N}$.
(c) Because $2800 \mathrm{~N}<6000 \mathrm{~N}$, the ball will not fracture your forehead, but will fracture your cheekbone because $2800 \mathrm{~N}>1300 \mathrm{~N}$.
Assess: A 90 mph fastball travels at $(90 \mathrm{mph})(1609.3 \mathrm{~m} / \mathrm{mile})(1 \mathrm{~h} / 3600 \mathrm{~s})=40 \mathrm{~m} / \mathrm{s}$, so it will not fracture your forehead, but it will fracture your cheekbone. This explains why baseball helmets protect the cheekbone.
7.29. Model: The rock ( R ) and Bob (b) are the two objects in our system, and will be treated as particles. We will also use the constant-acceleration kinematic equations.

## Visualize:

## Pictorial representation



Solve: (a) Bob exerts a forward force $\vec{F}_{\mathrm{B} \text { on } \mathrm{R}}$ on the rock to accelerate it forward. The rock's acceleration is calculated as follows:

$$
v_{1 \mathrm{R}}^{2}=v_{0 \mathrm{R}}^{2}+2 a_{0 \mathrm{R}} \Delta x \Rightarrow a_{\mathrm{R}}=\frac{v_{1 \mathrm{R}}^{2}}{2 \Delta x}=\frac{(30 \mathrm{~m} / \mathrm{s})^{2}}{2(1.0 \mathrm{~m})}=450 \mathrm{~m} / \mathrm{s}^{2}
$$

The force is calculated from Newton's second law:

$$
F_{\mathrm{B} \text { on } \mathrm{R}}=m_{\mathrm{R}} a_{\mathrm{R}}=(0.500 \mathrm{~kg})\left(450 \mathrm{~m} / \mathrm{s}^{2}\right)=225 \mathrm{~N}
$$

Bob exerts a force of $2.3 \times 10^{2} \mathrm{~N}$ on the rock.
(b) Because Bob pushes on the rock, the rock pushes back on Bob with a force $\vec{F}_{\mathrm{R} \text { on } \mathrm{B}}$. Forces $\vec{F}_{\mathrm{R} \text { on B }}$ and $\vec{F}_{\mathrm{B} \text { on R }}$ are an action/reaction pair, so $F_{\mathrm{R} \text { on } \mathrm{B}}=F_{\mathrm{B} \text { on } \mathrm{R}}=225 \mathrm{~N}$. The force causes Bob to accelerate backward with an acceleration of

$$
a_{\mathrm{B}}=\frac{\left(F_{\text {net on B }}\right)_{x}}{m_{\mathrm{B}}}=-\frac{F_{\mathrm{R} \text { on } \mathrm{B}}}{m_{\mathrm{B}}}=-\frac{225 \mathrm{~N}}{75 \mathrm{~kg}}=-3.0 \mathrm{~m} / \mathrm{s}^{2}
$$

This is a rather large acceleration, but it lasts only until Bob releases the rock. We can determine the time interval by returning to the kinematics of the rock:

$$
v_{1 \mathrm{R}}=v_{0 \mathrm{R}}+a_{\mathrm{R}} \Delta t=a_{\mathrm{R}} \Delta t \Rightarrow \Delta t=\frac{v_{1 \mathrm{R}}}{a_{\mathrm{R}}}=0.0667 \mathrm{~s}
$$

At the end of this interval, Bob's velocity is

$$
v_{1 \mathrm{~B}}=v_{0 \mathrm{~B}}+a_{\mathrm{B}} \Delta t=a_{\mathrm{B}} \Delta t=-0.20 \mathrm{~m} / \mathrm{s}
$$

Thus his recoil speed is $0.20 \mathrm{~m} / \mathrm{s}$.
7.30. Model: The boy $(B)$ and the crate $(C)$ are the two objects in our system, and they will be treated in the particle model. We will also use the static and kinetic friction models.

## Visualize:

## Pictorial representation



Solve: The fact that the boy's feet occasionally slip means that the maximum force of static friction must exist between the boy's feet and the sidewalk. That is, $f_{\mathrm{sB}}=\mu_{\mathrm{sB}} n_{\mathrm{B}}$. Also $f_{\mathrm{kC}}=\mu_{\mathrm{kC}} n_{\mathrm{C}}$.
Newton's second law applied to the crate gives

$$
\begin{gathered}
\sum\left(F_{\mathrm{on} \mathrm{C}}\right)_{y}=n_{\mathrm{C}}-\left(F_{\mathrm{G}}\right)_{\mathrm{C}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{C}}=m_{\mathrm{C}} g \\
\sum\left(F_{\mathrm{on} \mathrm{C}}\right)_{x}=F_{\mathrm{B} \text { on } \mathrm{C}}-f_{\mathrm{kC}}=0 \mathrm{~N} \Rightarrow \quad F_{\mathrm{B} \text { on } \mathrm{C}}=f_{\mathrm{kC}}=\mu_{\mathrm{kC}} n_{\mathrm{C}}=\mu_{\mathrm{kC}} m_{\mathrm{C}} g
\end{gathered}
$$

Newton's second law for the boy is

$$
\begin{gathered}
\sum\left(F_{\text {on B }}\right)_{y}=n_{\mathrm{B}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{B}}=m_{\mathrm{B}} g \\
\sum\left(F_{\mathrm{on} \mathrm{~B}}\right)_{x}=f_{\mathrm{sB}}-F_{\mathrm{C} \text { on } \mathrm{B}}=0 \mathrm{~N} \Rightarrow F_{\mathrm{C} \text { on } \mathrm{B}}=f_{\mathrm{sB}}=\mu_{\mathrm{sB}} n_{\mathrm{B}}=\mu_{\mathrm{sB}} m_{\mathrm{B}} g
\end{gathered}
$$

$\vec{F}_{\mathrm{C} \text { on } \mathrm{B}}$ and $\vec{F}_{\mathrm{B} \text { on } \mathrm{C}}$ are an action/reaction pair, so

$$
F_{\mathrm{C} \text { on } \mathrm{B}}=F_{\mathrm{B} \text { on } \mathrm{C}} \Rightarrow \mu_{\mathrm{sB}} m_{\mathrm{B}} g=\mu_{\mathrm{kC}} m_{\mathrm{C}} g \quad \Rightarrow \quad m_{\mathrm{C}}=\frac{\mu_{\mathrm{SB}} m_{\mathrm{B}}}{\mu_{\mathrm{kC}}}=\frac{(0.8)(50 \mathrm{~kg})}{(0.2)}=2 \times 10^{2} \mathrm{~kg}
$$

7.31. Model: Assume package A and package B are particles. Use the model of kinetic friction and the constantacceleration kinematic equations.

## Visualize:

Pictorial representation


| Known |  |
| :--- | :--- |
| $m_{\mathrm{A}}=5.0 \mathrm{~kg}$ | $m_{\mathrm{B}}=10 \mathrm{~kg}$ |
| $\theta=20^{\circ}$ | $\mu_{\mathrm{kA}}=0.20$ |
| $\mu_{\mathrm{kB}}=0.15$ |  |
| $x_{0}=v_{0 x}=t_{0}=0$ | $x_{1}=2.0 \mathrm{~m}$ |

Find
$t_{1}$


Package B


Package A

Solve: Package B has a smaller coefficient of friction, so its acceleration down the ramp is greater than that of package A. It will therefore push against package A and, by Newton's third law, package A will push back on B. The acceleration constraint is $a_{\mathrm{A}}=a_{\mathrm{B}} \equiv a$.
Newton's second law applied to each package gives

$$
\begin{aligned}
& \sum\left(F_{\text {on } \mathrm{A}}\right)_{x}=F_{\mathrm{B} \text { on } \mathrm{A}}+\left(F_{\mathrm{G}}\right)_{\mathrm{A}} \sin \theta-f_{\mathrm{kA}}=m_{\mathrm{A}} a \\
& F_{\mathrm{B} \text { on } \mathrm{A}}+m_{\mathrm{A}} g \sin \theta-\mu_{\mathrm{kA}}\left(m_{\mathrm{A}} g \cos \theta\right)=m_{\mathrm{A}} a \\
& \sum\left(F_{\text {on } \mathrm{B}}\right)_{x}=-F_{\mathrm{A} \text { on } \mathrm{B}}-f_{\mathrm{kB}}+\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \sin \theta=m_{\mathrm{B}} a \\
& -F_{\mathrm{A} \text { on } \mathrm{B}}-\mu_{\mathrm{kB}}\left(m_{\mathrm{B}} g \cos \theta\right)+m_{\mathrm{B}} g \sin \theta=m_{\mathrm{B}} a
\end{aligned}
$$

where we have used $n_{\mathrm{A}}=m_{\mathrm{A}} \cos \theta g$ and $n_{\mathrm{B}}=m_{\mathrm{B}} \cos \theta g$. Adding the two force equations, and using $F_{\mathrm{A} \text { on } \mathrm{B}}=F_{\mathrm{B} \text { on A }}$ because they are an action/reaction pair, we get

$$
a=g \sin \theta-\frac{\left(\mu_{\mathrm{kA}} m_{\mathrm{A}}+\mu_{\mathrm{kB}} m_{\mathrm{B}}\right)(g \cos \theta)}{m_{\mathrm{A}}+m_{\mathrm{B}}}=\frac{[(020)(5.0 \mathrm{~kg})+(0.15)(10 \mathrm{~kg})]\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos \left(20^{\circ}\right)}{5.0 \mathrm{~kg}+10 \mathrm{~kg}}=1.82 \mathrm{~m} / \mathrm{s}^{2}
$$

Finally, using $x_{1}=x_{0}+v_{0}\left(t_{1}-t_{0}\right)+\frac{1}{2} a\left(t_{1}-t_{0}\right)^{2}$, we find

$$
2.0 \mathrm{~m}=0 \mathrm{~m}+0 \mathrm{~m}+\frac{1}{2}\left(1.82 \mathrm{~m} / \mathrm{s}^{2}\right)\left(t_{1}-0 \mathrm{~s}\right)^{2} \Rightarrow t_{1}=\sqrt{2(2.0 \mathrm{~m}) /\left(1.82 \mathrm{~m} / \mathrm{s}^{2}\right)}=1.5 \mathrm{~s}
$$

7.32. Model: The two blocks form a system of interacting objects. We shall treat them as particles.

Visualize: Please refer to Figure P7.32.


Solve: It is possible that the left-hand block (Block L) is accelerating down the slope faster than the right-hand block (Block R), causing the string to be slack (zero tension). If that were the case, we would get a zero or negative answer for the tension in the string. Newton's first law applied in the $y$-direction on Block L yields

$$
\left(\sum F_{\mathrm{L}}\right)_{y}=0=n_{\mathrm{L}}-\left(F_{\mathrm{G}}\right)_{\mathrm{L}} \cos \left(20^{\circ}\right) \Rightarrow n_{\mathrm{L}}=m_{\mathrm{L}} g \cos \left(20^{\circ}\right)
$$

Therefore

$$
\left(f_{\mathrm{k}}\right)_{\mathrm{L}}=\left(\mu_{\mathrm{k}}\right)_{\mathrm{L}} m_{\mathrm{L}} g \cos \left(20^{\circ}\right)=(0.20)(1.0 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right) \cos \left(20^{\circ}\right)=1.84 \mathrm{~N}
$$

A similar analysis of the forces in the $y$-direction on Block R gives $\left(f_{\mathrm{k}}\right)_{\mathrm{R}}=1.84 \mathrm{~N}$ as well. Using Newton's second law in the $x$-direction for Block L gives

$$
\left(\sum F_{\mathrm{L}}\right)_{x}=m_{\mathrm{L}} a=T_{\mathrm{R} \text { on } \mathrm{L}}-\left(f_{\mathrm{k}}\right)_{\mathrm{L}}+\left(F_{\mathrm{G}}\right)_{\mathrm{L}} \sin \left(20^{\circ}\right) \quad \Rightarrow \quad m_{\mathrm{L}} a=T_{\mathrm{R} \text { on } \mathrm{L}}-1.84 \mathrm{~N}+m_{\mathrm{L}} g \sin \left(20^{\circ}\right)
$$

For Block R,

$$
\left(\sum F_{\mathrm{R}}\right)_{x}=m_{\mathrm{R}} a=\left(F_{\mathrm{G}}\right)_{\mathrm{R}} \sin \left(20^{\circ}\right)-1.84 \mathrm{~N}-T_{\mathrm{L} \text { on } \mathrm{R}} \quad \Rightarrow \quad m_{\mathrm{R}} a=m_{\mathrm{R}} g \sin \left(20^{\circ}\right)-1.84 \mathrm{~N}-T_{\mathrm{L} \text { on } \mathrm{R}}
$$

Solving these two equations in the two unknowns $a$ and $T_{\mathrm{L} \text { on } \mathrm{R}}=T_{\mathrm{R} \text { on } \mathrm{L}} \equiv T$, we obtain $a=2.12 \mathrm{~m} / \mathrm{s}^{2}$ and $T=0.61 \mathrm{~N}$.
Assess: The tension in the string is positive, and is about $1 / 3$ of the kinetic friction force on each of the blocks, which is reasonable.
7.33. Model: The two blocks (1 and 2) form the system of interest and will be treated as particles. The ropes are assumed to be massless, and the model of kinetic friction will be used.

## Visualize:

## Pictorial representation



$$
\begin{aligned}
& \frac{\text { Known }}{T_{\text {pull }}=20 \mathrm{~N}} \\
& \mu_{\mathrm{k}}=0.40 \\
& m_{1}=1.0 \mathrm{~kg} \\
& m_{2}=2.0 \mathrm{~kg} \\
& \text { Find } \\
& \hline T_{\text {rope }} a
\end{aligned}
$$



Solve: (a) The separate free-body diagrams for the two blocks show that there are two action/reaction pairs. Notice how block 1 both pushes down on block 2(force $\vec{n}_{1}^{\prime}$ )and exerts a retarding friction force $\vec{f}_{2}$ top on the top surface of block 2 . Block 1 is in static equilibrium $\left(a_{1}=0 \mathrm{~m} / \mathrm{s}^{2}\right)$ but block 2 is accelerating to the right. Newton's second law for block 1 is

$$
\begin{aligned}
& \left(F_{\text {net on } 1}\right)_{x}=f_{1}-T_{\text {rope }}=0 \mathrm{~N} \quad \Rightarrow \quad T_{\text {rope }}=f_{1} \\
& \left(F_{\text {net on } 1}\right)_{y}=n_{1}-m_{1} g=0 \mathrm{~N} \quad \Rightarrow \quad n_{1}=m_{1} g
\end{aligned}
$$

Although block 1 is stationary, there is a kinetic force of friction because there is motion between block 1 and block 2 . The friction model means $f_{1}=\mu_{k} n_{1}=\mu_{k} m_{1} g$. Substitute this result into the $x$-equation to get the tension in the rope:

$$
T_{\text {rope }}=f_{1}=\mu_{\mathrm{k}} m_{1} g=(0.40)(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=3.9 \mathrm{~N}
$$

(b) Newton's second law for block 2 is

$$
\begin{aligned}
& a_{x} \equiv a=\frac{\left(F_{\text {net on 2 }}\right)_{x}}{m_{2}}=\frac{T_{\text {pull }}-f_{2 \text { top }}-f_{2 \text { bot }}}{m_{2}} \\
& a_{y}=0 \mathrm{~m} / \mathrm{s}^{2}=\frac{\left(F_{\text {net on } 2)_{y}}\right.}{m_{2}}=\frac{n_{2}-n_{1}^{\prime}-m_{2} g}{m_{2}}
\end{aligned}
$$

Forces $\vec{n}_{1}$ and $\vec{n}_{1}^{\prime}$ are an action/reaction pair, so $n_{1}^{\prime}=n_{1}=m_{1} g$. Substituting into the $y$-equation gives $n_{2}=\left(m_{1}+m_{2}\right) g$. This is not surprising because the combined weight of both objects presses down on the surface. The kinetic friction on the bottom surface of block 2 is then

$$
f_{2 \text { bot }}=\mu_{\mathrm{k}} n_{2}=\mu_{\mathrm{k}}\left(m_{1}+m_{2}\right) g
$$

The forces $\vec{f}_{1}$ and $\vec{f}_{2 \text { top }}$ are an action/reaction pair, so $f_{2 \text { bot }}=f_{1}=\mu_{\mathrm{k}} m_{1} g$. Inserting these friction results into the $x$-equation gives

$$
\begin{aligned}
a & =\frac{\left(F_{\text {net on 2 }}\right)_{x}}{m_{2}}=\frac{T_{\text {pull }}-\mu_{\mathrm{k}} m_{1} g-\mu_{\mathrm{k}}\left(m_{1}+m_{2}\right) g}{m_{2}} \\
& =\frac{20 \mathrm{~N}-(0.40)(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)-(0.40)(1.0 \mathrm{~kg}+2.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{2.0 \mathrm{~kg}}=2.2 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

7.34. Model: The $3-\mathrm{kg}$ and $4-\mathrm{kg}$ blocks constitute the system and are to be treated as particles. The models of kinetic and static friction and the constant-acceleration kinematic equations will be used.

## Visualize:

## Pictorial representation

Known
$m_{3}=3.0 \mathrm{~kg}$
$m_{4}=4.0 \mathrm{~kg}$
$\mu_{\mathrm{s}}($ Block on block $)=0.60$
$\mu_{\mathrm{k}}($ Block on floor $)=0.20$
$x_{0}=v_{0 x}=t_{0}=0$
$x_{1}=5.0 \mathrm{~m}$

Block 4
Find
$t_{1}$ without sliding

Solve: The minimum time will be achieved when static friction is at its maximum possible value. Newton's second law for the 4-kg block is

$$
\begin{gathered}
\sum\left(F_{\text {on } 4}\right)_{y}=n_{3 \text { on } 4}-\left(F_{\mathrm{G}}\right)_{4}=0 \mathrm{~N} \Rightarrow n_{3 \text { on } 4}=\left(F_{\mathrm{G}}\right)_{4}=m_{4} g=(4.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=39.2 \mathrm{~N} \\
f_{\mathrm{s} 4}=\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n_{3 \text { on } 4}=(0.60)(39.2 \mathrm{~N})=23.5 \mathrm{~N}
\end{gathered}
$$

Newton's second law for the $3-\mathrm{kg}$ block is

$$
\sum\left(F_{\text {on } 3}\right)_{y}=n_{3}-n_{4 \text { on } 3}-\left(F_{\mathrm{G}}\right)_{3}=0 \mathrm{~N} \Rightarrow n_{3}=n_{4 \text { on } 3}+\left(F_{\mathrm{G}}\right)_{3}=39.2 \mathrm{~N}+(3.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=68.6 \mathrm{~N}
$$

Friction forces $f$ and $f_{\mathrm{s} 4}$ are an action/reaction pair. Thus

$$
\begin{gathered}
\sum\left(F_{\text {on } 3}\right)_{x}=f_{\mathrm{s} 3}-f_{\mathrm{k} 3}=m_{3} a_{3} \Rightarrow f_{\mathrm{s} 4}-\mu_{\mathrm{k}} n_{3}=m_{3} a_{3} \Rightarrow 23.5 \mathrm{~N}-(0.20)(68.6 \mathrm{~N})=(3.0 \mathrm{~kg}) a_{3} \\
a_{3}=3.27 \mathrm{~m} / \mathrm{s}^{2}
\end{gathered}
$$

Since block 3 does not slip, this is also the acceleration of block 4. The time is calculated as follows:

$$
x_{1}-x_{0}+v_{0 x}\left(t_{1}-t_{0}\right)+\frac{1}{2} a\left(t_{1}-t_{0}\right)^{2} \Rightarrow 5.0 \mathrm{~m}=0 \mathrm{~m}+0 \mathrm{~m}+\frac{1}{2}\left(3.27 \mathrm{~m} / \mathrm{s}^{2}\right)\left(t_{1}-0 \mathrm{~s}\right)^{2} \Rightarrow t_{1}=1.8 \mathrm{~s}
$$

7.35. Model: Blocks 1 and 2 make up the system of interest and will be treated as particles. Assume a massless rope and frictionless pulley.
Visualize:

## Pictorial representation



Solve: The blocks accelerate with the same magnitude but in opposite directions. Thus the acceleration constraint is $a_{2}=a=-a_{1}$, where a will have a positive value. There are two real action/reaction pairs. The two tension forces will act as if they are action/reaction pairs because we are assuming a massless rope and a frictionless pulley. Make sure you understand why the friction forces point in the directions shown in the free-body diagrams, especially force $\vec{f}_{1}^{\prime}$ exerted on block 2 by block 1. We have quite a few pieces of information to include. First, Newton's second law applied to blocks 1 and 2 gives

$$
\begin{aligned}
& \left(\vec{F}_{\text {net on } 1}\right)_{x}=f_{1}-T_{1}=\mu_{\mathrm{k}} n_{1}-T_{1}=m_{1} a_{1}=-m_{1} a \\
& \left(F_{\text {net on } 1}\right)_{y}=n_{1}-m_{1} g=0 \mathrm{~N} \Rightarrow n_{1}=m_{1} g \\
\left(F_{\text {net on } 2}\right)_{x}= & T_{\text {pull }}-f_{1}^{\prime}-f_{2}-T_{2}=T_{\text {pull }}-f_{1}^{\prime}-\mu_{\mathrm{k}} n_{2}-T_{2}=m_{2} a_{2}=m_{2} a \\
\left(F_{\text {net on 2 } 2}\right)_{y}= & n_{2}-n_{1}^{\prime}-m_{2} g=0 \mathrm{~N} \Rightarrow n_{2}=n_{1}^{\prime}+m_{2} g
\end{aligned}
$$

We've already used the kinetic friction model in both $x$-equations. Next, Newton's third law gives

$$
n_{1}^{\prime}=n_{1}=m_{1} g \quad f_{1}^{\prime}=f_{1}=\mu_{\mathrm{k}} n_{1}=\mu_{\mathrm{k}} m_{1} g \quad T_{1}=T_{2}=T
$$

Knowing $n_{1}^{\prime}$, we can now use the $y$-equation of block 2 to find $n_{2}$. Substitute all these pieces into the two $x$-equations, and we end up with two equations with two unknowns:

$$
\mu_{\mathrm{k}} m_{1} g-T=-m_{1} a \quad T_{\text {pull }}-T-\mu_{\mathrm{k}} m_{1} g-\mu_{\mathrm{k}}\left(m_{1}+m_{2}\right) g=m_{2} a
$$

Subtract the first equation from the second to get

$$
\begin{aligned}
& T_{\text {pull }}-\mu_{\mathrm{k}}\left(3 m_{1}+m_{2}\right) g=\left(m_{1}+m_{2}\right) a \\
& a=\frac{T_{\text {pull }}-\mu_{\mathrm{k}}\left(3 m_{1}+m_{2}\right) g}{m_{1}+m_{2}}=\frac{20 \mathrm{~N}-(0.30)[3(1.0 \mathrm{~kg})+2.0 \mathrm{~kg}]\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{1.0 \mathrm{~kg}+2.0 \mathrm{~kg}}=1.8 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

7.36. Model: Blocks 1 and 2 make up the system of interest and will be treated as particles. We shall use the kinetic friction model.

## Visualize:



Block A


Block B

Notice that the coordinate system of for block B is rotated so that the motion in the positive $x$-direction is consistent between the two free-body diagrams.
Solve: The blocks are constrained to have the same magnitude acceleration. Applying Newton's second law to block B gives

$$
\sum(F)_{y}=-T+\left(F_{G}\right)_{\mathrm{B}}=m a \Rightarrow T-m g=-m a
$$

Applying Newton's second law in both the $x$-and $y$-directions to the block A gives

$$
\begin{aligned}
& \sum(F)_{y}=n-\left(F_{G}\right)_{\mathrm{A}}=0 \Rightarrow n=M g \\
& \sum(F)_{x}=T=M a \Rightarrow T=M a
\end{aligned}
$$

Using the first equation to eliminate the acceleration $a$ gives the tension:

$$
T=M a=M(g-T / m) \quad \Rightarrow \quad T=\frac{m M g}{m+M}
$$

Assess: The result is positive, as it should be for our choice of coordinate system. Consider $m=0$. In this case, $T=0$, as expected. For $m \quad M$, the tension is independent of the mass of the hanging block because its acceleration will be $g$, as we can see by solving for the acceleration:

$$
a=-\frac{T}{m}+g=g-\frac{M g}{m+M} \rightarrow g \text { for } m \quad M
$$

7.37. Model: The sled (S) and the box (B) will be treated in the particle model, and the model of friction will be used. Refer to Table 6.1 for the required friction coefficients.

## Visualize:

## Pictorial representation


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In the sled's free-body diagram $n_{\mathrm{S}}$ is the normal (contact) force on the sled due to the snow. Similarly $f_{\mathrm{kS}}$ is the force of kinetic friction on the sled due to snow.
Solve: Newton's second law on the box in the $y$-direction is

$$
n_{\mathrm{S} \text { on } \mathrm{B}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \cos \left(20^{\circ}\right)=0 \mathrm{~N} \Rightarrow n_{\mathrm{S} \text { on } \mathrm{B}}=(10 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos \left(20^{\circ}\right)=92.1 \mathrm{~N}
$$

The static friction force $\vec{f}_{\mathrm{S}}$ on B accelerates the box. The maximum acceleration occurs when static friction reaches its maximum possible value.

$$
\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{S}} n_{\mathrm{S} \text { on } \mathrm{B}}=(0.50)(92.1 \mathrm{~N})=46.1 \mathrm{~N}
$$

Newton's second law along the $x$-direction thus gives the maximum acceleration

$$
f_{\mathrm{S} \text { on } \mathrm{B}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \sin \left(20^{\circ}\right)=m_{\mathrm{B}} a \Rightarrow 46.1 \mathrm{~N}-(10 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin \left(20^{\circ}\right)=(10 \mathrm{~kg}) a \Rightarrow a=1.25 \mathrm{~m} / \mathrm{s}^{2}
$$

Newton's second law for the sled along the $y$-direction is

$$
\begin{gathered}
n_{\mathrm{S}}-n_{\mathrm{B} \text { on } \mathrm{S}}-\left(F_{\mathrm{G}}\right)_{\mathrm{S}} \cos \left(20^{\circ}\right)=0 \mathrm{~N} \\
n_{\mathrm{S}}=n_{\mathrm{B} \text { on } \mathrm{S}}+m_{\mathrm{S}} g \cos \left(20^{\circ}\right)=(92.1 \mathrm{~N})+(20 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos \left(20^{\circ}\right)=276.3 \mathrm{~N}
\end{gathered}
$$

Therefore, the force of friction on the sled by the snow is

$$
f_{\mathrm{kS}}=\left(\mu_{\mathrm{k}}\right) n_{\mathrm{S}}=(0.06)(276.3 \mathrm{~N})=16.6 \mathrm{~N}
$$

Newton's second law along the $x$-direction is

$$
T_{\text {pull }}-w_{\mathrm{S}} \sin \left(20^{\circ}\right)-f_{\mathrm{kS}}-f_{\mathrm{B} \text { on } \mathrm{S}}=m_{\mathrm{S}} a
$$

The friction force $f_{\mathrm{B} \text { on } \mathrm{S}}=f_{\mathrm{S} \text { on } \mathrm{B}}$ because these are an action/reaction pair. We're using the maximum acceleration, so the maximum tension is

$$
T_{\max }-(20 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \sin \left(20^{\circ}\right)-16.6 \mathrm{~N}-46.1 \mathrm{~N}=(20 \mathrm{~kg})\left(1.25 \mathrm{~m} / \mathrm{s}^{2}\right)=160 \mathrm{~N}
$$

7.38. Model: The masses $m$ and $M$ are to be treated in the particle model. We will also assume a massless rope and frictionless pulley, and use the constant-acceleration kinematic equations for $m$ and $M$.
Visualize:

## Pictorial representation



Solve: Using $y_{1}=y_{0}+v_{0 y}\left(t_{1}-t_{0}\right)+\frac{1}{2} a_{\mathrm{M}}\left(t_{1}-t_{0}\right)^{2}$,

$$
(-1.0 \mathrm{~m})=0 \mathrm{~m}+0 \mathrm{~m}+\frac{1}{2} a_{\mathrm{M}}(6.0 \mathrm{~s}-0 \mathrm{~s})^{2} \Rightarrow a_{\mathrm{M}}=-0.0556 \mathrm{~m} / \mathrm{s}^{2}
$$

Newton's second law for $m$ and $M$ gives

$$
\sum\left(F_{\text {on m }}\right)_{y}=T_{\mathrm{R} \text { on } \mathrm{m}}-\left(F_{\mathrm{G}}\right)_{\mathrm{m}}=m a_{\mathrm{m}} \quad \sum\left(F_{\text {on } \mathrm{M}}\right)_{y}=T_{\mathrm{R} \text { on } \mathrm{M}}-\left(F_{\mathrm{G}}\right)_{\mathrm{M}}=M a_{\mathrm{M}}
$$

The acceleration constraint is $a_{m}=-a_{M}$. Also, the tensions are an pseudo-action/reaction pair, so $T_{\mathrm{R} \text { on } \mathrm{m}}=T_{\mathrm{R} \text { on } \mathrm{M}}$. With these, the second-law equations become

$$
\begin{gathered}
T_{\mathrm{R} \text { on } \mathrm{M}}-M g=M a_{\mathrm{M}} \\
T_{\mathrm{R} \text { on } \mathrm{M}}-m g=-m a_{\mathrm{M}}
\end{gathered}
$$

Subtracting the second from the first gives

$$
\begin{aligned}
-M g+m g & =M a_{\mathrm{M}}+m a_{\mathrm{M}} \\
m & =M\left[\frac{g+a_{\mathrm{M}}}{g-a_{\mathrm{M}}}\right] \\
& =(100 \mathrm{~kg})\left[\frac{9.8 \mathrm{~m} / \mathrm{s}^{2}-0.556 \mathrm{~m} / \mathrm{s}^{2}}{9.8 \mathrm{~m} / \mathrm{s}^{2}+0.556 \mathrm{~m} / \mathrm{s}^{2}}\right]=99 \mathrm{~kg}
\end{aligned}
$$

Assess: Note that $a_{\mathrm{m}}=-a_{\mathrm{M}}=0.0556 \mathrm{~m} / \mathrm{s}^{2}$. For such a small acceleration, the $1 \%$ mass difference seems reasonable.
7.39. Model: Use the particle model for the block of mass $M$ and the two massless pulleys. Additionally, the rope is massless and the pulleys are frictionless. The block is kept in place by an applied force $\vec{F}$.

## Visualize:

## Pictorial representation



Solve: Since there is no friction on the pulleys, $T_{2}=T_{3}=T_{5}$. Newton's second law for mass $M$ gives

$$
T_{1}-F_{\mathrm{G}}=0 \mathrm{~N} \Rightarrow T_{1}=M g=(10.2 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=100 \mathrm{~N}
$$

Newton's second law for the small pulley is

$$
T_{2}+T_{3}-T_{1}=0 \mathrm{~N} \quad \Rightarrow \quad T_{2}=T_{3}=\frac{T_{1}}{2}=50 \mathrm{~N}=T_{5}=F
$$

Newton's second law for the large pulley is

$$
T_{4}-T_{2}-T_{3}-T_{5}=0 \mathrm{~N} \Rightarrow T_{4}=T_{2}+T_{3}+T_{5}=150 \mathrm{~N}
$$

7.40. Model: Assume the particle model for $m_{1}, m_{2}$, and $m_{3}$, and the model of kinetic friction. Assume the ropes to be massless, and the pulleys to be frictionless and massless.

## Visualize:

## Pictorial representation



Solve: Newton's second law for $m_{1}$ gives $T_{1}-\left(F_{\mathrm{G}}\right)_{1}=m_{1} a_{1}$. Newton's second law for $m_{2}$ gives

$$
\begin{aligned}
& \sum\left(F_{\mathrm{on} m_{2}}\right)_{y}=n_{2}-\left(F_{\mathrm{G}}\right)_{2}=0 \mathrm{~N} \Rightarrow n_{2}=m_{2} g \\
& \sum\left(F_{\text {on } m_{2}}\right)_{x}=T_{2}-f_{\mathrm{k} 2}-T=m_{2} a_{2} \Rightarrow T_{2}-\mu_{\mathrm{k}} n_{2}-T_{1}=m_{2} a_{2}
\end{aligned}
$$

Newton's second law for $m_{3}$ gives $T_{2}-\left(F_{\mathrm{G}}\right)_{3}=m_{3} a_{3}$. Since $m_{1}, m_{2}$, and $m_{3}$ move together, $a_{1}=a_{2}=-a_{3} \equiv a$. The equations for the three masses thus become

$$
T_{1}-\left(F_{\mathrm{G}}\right)_{1}=m_{1} a \quad T_{2}-\mu_{\mathrm{k}} n_{2}-T_{1}=m_{2} a \quad T_{2}-\left(F_{\mathrm{G}}\right)_{3}=-m_{3} a
$$

Subtracting the third equation from the sum of the first two equations yields:

$$
\begin{gathered}
-\left(F_{\mathrm{G}}\right)_{1}-\mu_{\mathrm{k}} n_{2}+\left(F_{\mathrm{G}}\right)_{3}=-m_{1} g-\mu_{\mathrm{k}} m_{2} g+m_{3} g=\left(m_{1}+m_{2}+m_{3}\right) a \\
a=\frac{-m_{1} g-\mu_{\mathrm{k}} m_{2} g+m_{3} g}{\left(m_{1}+m_{2}+m_{3}\right)}=\frac{-1.0 \mathrm{~kg}-(0.30)(2.0 \mathrm{~kg})+3.0 \mathrm{~kg}}{1.0 \mathrm{~kg}+2.0 \mathrm{~kg}+3.0 \mathrm{~kg}}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=2.3 \mathrm{~m} / \mathrm{s}^{2}
\end{gathered}
$$

7.41. Model: Assume the particle model for the two blocks, and the model of kinetic and static friction. Visualize:

## Pictorial representation



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Solve: (a) If the mass $m$ is too small, the hanging 2.0 kg mass will pull it up the slope. We want to find the smallest mass that will stick as a result of friction. The smallest mass will be the one for which the force of static friction is at its maximum possible value: $f_{s}=\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n$. As long as the mass $m$ is stuck, both blocks are at rest with $\vec{F}_{\text {net }}=0 \mathrm{~N}$. In this situation, Newton's second law for the hanging mass $M$ gives

$$
\left(F_{\text {net }}\right)_{x}=-T_{\mathrm{M}}+M g=0 \mathrm{~N} \Rightarrow T_{\mathrm{M}}=M g=(2.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=19.6 \mathrm{~N}
$$

For the smaller mass $m$,

$$
\left(F_{\text {net }}\right)_{x}=T_{\mathrm{m}}-f_{\mathrm{s}}-m g \sin \theta=0 \mathrm{~N} \quad\left(F_{\mathrm{net}}\right)_{y}=n-m g \cos \theta \quad \Rightarrow \quad n=m g \cos \theta
$$

For a massless string and frictionless pulley, forces $\vec{T}_{\mathrm{m}}$ and $\vec{T}_{\mathrm{M}}$ act as if they are an action/reaction pair. Thus $T_{\mathrm{m}}=T_{\mathrm{M}}$. Mass $m$ is a minimum when $f_{\mathrm{s}}=\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n=\mu_{\mathrm{s}} m g \cos \theta$. Substituting these expressions into the $x$-equation for m gives

$$
\begin{aligned}
& T_{\mathrm{M}}-\mu_{\mathrm{s}} m g \cos \theta-m g \sin \theta=0 \mathrm{~N} \\
& m=\frac{T_{\mathrm{M}}}{\left(\mu_{\mathrm{s}} \cos \theta+\sin \theta\right) g}=\frac{19.6 \mathrm{~N}}{\left[(0.80) \cos \left(20^{\circ}\right)+\sin \left(20^{\circ}\right)\right]\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}=1.83 \mathrm{~kg}
\end{aligned}
$$

or 1.8 kg to two significant figures.
(b) Because $\mu_{\mathrm{k}}<\mu_{\mathrm{S}}$ the 1.8 kg block will begin to slide up the ramp and the 2.0 kg mass will begin to fall if the block is nudged ever so slightly. In this case, the net force and the acceleration are not zero. Notice how, in the pictorial representation, we chose different coordinate systems for the two masses. The magnitudes of the accelerations are the same because the blocks are tied together. Thus, the acceleration constraint is $a_{\mathrm{m}}=a_{\mathrm{M}} \equiv a$, where $a$ will have a positive value. Newton's second law for block M gives

$$
\left(F_{\text {net }}\right)_{x}=-T+M g=M a_{\mathrm{M}}=M a
$$

For block m we have

$$
\left(F_{\text {net }}\right)_{x}=T-f_{\mathrm{k}}-m g \sin \theta=T-\mu_{\mathrm{k}} m g \cos \theta-m g \sin \theta=m a_{\mathrm{m}}=m a
$$

In writing these equations, we used Newton's third law to obtain $T_{\mathrm{m}}=T_{\mathrm{M}}=T$. Also, notice that the $x$-equation and the friction model for block m don't change, except for $\mu_{\mathrm{s}}$ becoming $\mu_{\mathrm{k}}$, so we already know the expression for $f_{\mathrm{k}}$ from part (a). Notice that the tension in the string is not the gravitational force $M g$. We have two equations with the two unknowns $T$ and $a$ :

$$
M g-T=M a \quad T-\left(\mu_{\mathrm{k}} \cos \theta+\sin \theta\right) m g=m a
$$

Adding the two equations to eliminate $T$ gives

$$
\begin{aligned}
M g-\left(\mu_{\mathrm{k}} \cos \theta+\sin \theta\right) m g & =M a+m a \\
a & =g \frac{M-\left(\mu_{\mathrm{k}} \cos \theta+\sin \theta\right) m}{M+m} \\
& =\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \frac{2.0 \mathrm{~kg}-\left[(0.50) \cos \left(20^{\circ}\right)+\sin \left(20^{\circ}\right)\right](1.83 \mathrm{~kg})}{2.0 \mathrm{~kg}+1.83 \mathrm{~kg}} 1.3 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

7.42. Model: Assume the particle model for the two blocks and use the friction model.

Visualize:
Pictorial representation



Solve: (a) The slope is frictionless, so the blocks stay in place only if held. Once $m$ is released, the blocks will move one way or the other. As long as m is held, the blocks are in static equilibrium with $\vec{F}_{\text {net }}=0 \mathrm{~N}$. In this case, Newton's second law for the hanging block M is

$$
\left(F_{\text {net on } \mathrm{M}}\right)_{y}=T_{\mathrm{M}}-M g=0 \mathrm{~N} \Rightarrow T_{\mathrm{M}}=M g=19.6 \mathrm{~N}
$$

Because the string is massless and the pulley is frictionless, $T_{\mathrm{M}}=T_{\mathrm{m}}=T=20 \mathrm{~N}$ (to two significant figures).
(b) The free-body diagram shows box $m$ after it is released. Whether it moves up or down the slope depends on whether the acceleration $a$ is positive or negative. The acceleration constraint is $\left(a_{\mathrm{m}}\right)_{x}-\left(a_{\mathrm{M}}\right)_{y} \equiv a$ Newton's second law for each system gives

$$
\left(F_{\text {net on } \mathrm{m}}\right)_{x}=T-m g \sin \theta=m\left(a_{\mathrm{m}}\right)_{x}=m a \quad\left(F_{\text {net on M }}\right)_{y}=T-M g=M\left(a_{\mathrm{M}}\right)_{y}=-M a
$$

We have two equations in two unknowns. Subtract the second from the first to eliminate $T$ :

$$
-m g \sin \theta+M g=(m+M) a \Rightarrow a=\frac{M-m \sin \theta}{M+m} g=\frac{2.0 \mathrm{~kg}-(4.0 \mathrm{~kg}) \sin \left(35^{\circ}\right)}{2.0 \mathrm{~kg}+4.0 \mathrm{~kg}}=-0.48 \mathrm{~m} / \mathrm{s}^{2}
$$

Since $a<0 \mathrm{~m} / \mathrm{s}^{2}$, the box accelerates down the slope.
(c) It is now straightforward to compute $T=M g-M a=21 \mathrm{~N}$. Notice how the tension is larger than when the blocks were motionless.
7.43. Model: Use the particle model for the book (B) and the coffee cup (C), the models of kinetic and static friction, and the constant-acceleration kinematic equations.

## Visualize:



Solve: (a) Using $v_{1 x}^{2}=v_{0 x}^{2}+2 a\left(x_{1}-x_{0}\right)$, we find

$$
0 \mathrm{~m}^{2} / \mathrm{s}^{2}=(3.0 \mathrm{~m} / \mathrm{s})^{2}+2 a\left(x_{1}\right) \Rightarrow a x_{1}=-4.5 \mathrm{~m}^{2} / \mathrm{s}^{2}
$$

To find $x_{1}$, we must first find $a$. Newton's second law applied to the book and the coffee cup gives

$$
\begin{gathered}
\sum\left(F_{\text {on B }}\right)_{y}=n_{\mathrm{B}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \cos \left(20^{\circ}\right)=0 \mathrm{~N} \Rightarrow n_{\mathrm{B}}=(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right) \cos \left(20^{\circ}\right)=9.21 \mathrm{~N} \\
\sum\left(F_{\text {on B }}\right)_{x}=-T-f_{\mathrm{k}}-\left(F_{\mathrm{G}}\right)_{\mathrm{B}} \sin \left(20^{\circ}\right)=m_{\mathrm{B}} a_{\mathrm{B}} \quad \sum\left(F_{\text {on } \mathrm{C}}\right)_{y}=T-\left(F_{\mathrm{G}}\right)_{\mathrm{C}}=m_{\mathrm{C}} a_{\mathrm{C}}
\end{gathered}
$$

The last two equations can be rewritten, using $a_{\mathrm{C}}=a_{\mathrm{B}}=a$, as

$$
-T-\mu_{\mathrm{k}} n_{\mathrm{B}}-m_{\mathrm{B}} g \sin \left(20^{\circ}\right)=m_{\mathrm{B}} a \quad T-m_{\mathrm{C}} g=m_{\mathrm{C}} a
$$

Adding the two equations gives

$$
\begin{gathered}
a\left(m_{\mathrm{C}}+m_{\mathrm{B}}\right)=-g\left[m_{\mathrm{C}}+m_{\mathrm{B}} \sin \left(20^{\circ}\right)\right]-\mu_{\mathrm{k}}(9.21 \mathrm{~N}) \\
(1.5 \mathrm{~kg}) a=-\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)\left[0.500 \mathrm{~kg}+(1.0 \mathrm{~kg}) \sin 20^{\circ}\right]-(0.20)(9.21 \mathrm{~N}) \quad \Rightarrow \quad a=-6.73 \mathrm{~m} / \mathrm{s}^{2}
\end{gathered}
$$

Using this value for $a$, we can now find $x_{1}$ as follows:

$$
x_{1}=\frac{-4.5 \mathrm{~m}^{2} / \mathrm{s}^{2}}{a}=\frac{-4.5 \mathrm{~m}^{2} / \mathrm{s}^{2}}{-6.73 \mathrm{~m} / \mathrm{s}^{2}}=0.67 \mathrm{~m}
$$

(b) The maximum static friction force is $\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n_{\mathrm{B}}=(0.50)(9.21 \mathrm{~N})=4.60 \mathrm{~N}$. We'll see if the force $f_{\mathrm{s}}$ needed to keep the book in place is larger or smaller than $\left(f_{\mathrm{s}}\right)_{\max }$. When the cup is at rest, the string tension is $T=m_{\mathrm{C}} g$. Newton's first law for the book is

$$
\begin{aligned}
\sum\left(F_{\text {on } \mathrm{B}}\right)_{x} & =f_{\mathrm{s}}-T-w_{\mathrm{B}} \sin \left(20^{\circ}\right)=f_{\mathrm{s}}-m_{\mathrm{C}} g-m_{\mathrm{B}} g \sin \left(20^{\circ}\right)=0 \\
f_{\mathrm{s}} & =\left(M_{\mathrm{C}}+M_{\mathrm{B}} \sin 20^{\circ}\right) g=8.25 \mathrm{~N}
\end{aligned}
$$

Because $f_{\mathrm{s}}>\left(f_{\mathrm{s}}\right)_{\max }$, the book slides back down.
7.44. Model: Use the particle model for the cable car and the counterweight. Assume a massless cable.

Visualize:

## Pictorial representation

Known
$x_{0}=v_{0}=0$
$\theta_{1}=30^{\circ} \quad \theta_{2}=20^{\circ}$
$x_{1}=-200 \mathrm{~m} / \sin 30^{\circ}$
$=-400 \mathrm{~m}$
$\left(a_{1}\right)_{x}=\left(a_{2}\right)_{x}=a$
Find
$F_{\mathrm{B}} \quad v_{1}$


Solve: (a) Notice the separate coordinate systems for the cable car (object 1) and the counterweight (object 2). Forces $\vec{T}_{1}$ and $\vec{T}_{2}$ act as if they are an action/reaction pair. The braking force $\vec{F}_{\mathrm{B}}$ works with the cable tension $\vec{T}_{1}$ to allow the cable car to descend at a constant speed. Constant speed means dynamic equilibrium, so $\vec{F}_{\text {net }}=0 \mathrm{~N}$ for both systems. Newton's second law applied to the cable car gives

$$
\left(F_{\text {net on } 1}\right)_{x}=T_{1}+F_{\mathrm{B}}-m_{1} g \sin \theta_{1}=0 \mathrm{~N} \quad\left(F_{\text {net on } 1}\right)_{y}=n_{1}-m_{1} g \cos \theta_{1}=0 \mathrm{~N}
$$

Newton's second law applied to the counterweight gives

$$
\left(F_{\text {net on } 2}\right)_{x}=m_{2} g \sin \theta_{2}-T_{2}=0 \mathrm{~N} \quad\left(F_{\text {net on } 2}\right)_{y}=n_{2}-m_{2} g \cos \theta_{2}=0 \mathrm{~N}
$$

From the $x$-equation for the counterweight, $T_{2}=m_{2} g \sin \theta_{2}$. Because we can neglect the pulley's friction and the cable is assumed to be massless, $T_{1}=T_{2}$. Thus the $x$-equation for the cable car then becomes

$$
F_{\mathrm{B}}=m_{1} g \sin \theta_{1}-T_{1}=m_{1} g \sin \theta_{1}-m_{2} g \sin \theta_{2}=3770 \mathrm{~N}=3.8 \mathrm{kN}
$$

(b) If the brakes fail, then $F_{\mathrm{B}}=0 \mathrm{~N}$. The car will accelerate down the hill on one side while the counterweight accelerates up the hill on the other side. Both will have negative accelerations because of the direction of the acceleration vectors. The constraint is $a_{1 x}=a_{2 x}=a$, where $a$ will have a negative value. Using $T_{1}=T_{2}=T$, the two $x$-equations are

$$
\left(F_{\text {net on } 1}\right)_{x}=T-m_{1} g \sin \theta_{1}=m_{1} a_{1 x}=m_{1} a \quad\left(F_{\text {net on } 2}\right)_{x}=m_{2} g \sin \theta_{2}-T=m_{2} a_{2 x}=m_{2} a
$$

Note that the $y$-equations aren't needed in this problem. Add the two equations to eliminate $T$ :

$$
-m_{1} g \sin \theta_{1}+m_{2} g \sin \theta=\left(m_{1}+m_{2}\right) a \Rightarrow a=-\frac{m_{1} \sin \theta_{1}-m_{2} \sin \theta_{2}}{m_{1}+m_{2}} g=-0.991 \mathrm{~m} / \mathrm{s}^{2}
$$

Now we have a problem in kinematics. The speed at the bottom is calculated as follows:

$$
v_{1}^{2}=v_{0}^{2}+2 a\left(x_{1}-x_{0}\right)=2 a x_{1} \quad \Rightarrow \quad v_{1}=\sqrt{2 a x_{1}}=\sqrt{2\left(-0.991 \mathrm{~m} / \mathrm{s}^{2}\right)(-400 \mathrm{~m})}=28 \mathrm{~m} / \mathrm{s}
$$

Assess: A speed of approximately 60 mph as the cable car travels a distance of 2000 m along a frictionless slope of $30^{\circ}$ is reasonable.
7.45. Model: Assume the cable mass is negligible compared to the car mass and that the pulley is frictionless. Use the particle model for the two cars.
Visualize: Please refer to Figure P7.45.

## Pictorial representation



Solve: (a) The cars are moving at constant speed, so they are in dynamic equilibrium. Consider the descending car D. We can find the rolling friction force on car D, and then find the cable tension by applying Newton's first law. In the $y$-direction for car D ,

$$
\begin{aligned}
\left(F_{\text {net }}\right)_{y} & =0=n_{\mathrm{D}}-\left(F_{\mathrm{G}}\right)_{\mathrm{D}} \cos \left(35^{\circ}\right) \\
n_{\mathrm{D}} & =m_{\mathrm{D}} g \cos \left(35^{\circ}\right)
\end{aligned}
$$

So the rolling friction force on car D is

$$
\left(f_{\mathrm{R}}\right)_{\mathrm{D}}=\mu_{\mathrm{R}} n_{\mathrm{D}}=\mu_{\mathrm{R}} m_{\mathrm{D}} g \cos \left(35^{\circ}\right)
$$

Applying Newton's first law to car D in the $x$-direction gives

$$
\left(F_{n e t}\right)_{x}=T_{\mathrm{A} \text { on } \mathrm{D}}+\left(f_{\mathrm{R}}\right)_{\mathrm{D}}-\left(F_{\mathrm{G}}\right)_{\mathrm{D}} \sin \left(35^{\circ}\right)=0
$$

Thus,

$$
\begin{aligned}
T_{\mathrm{A} \text { on D }} & =m_{\mathrm{D}} g \sin \left(35^{\circ}\right)-\mu_{\mathrm{R}} m_{\mathrm{D}} g \cos \left(35^{\circ}\right) \\
& =(1500 \mathrm{~kg})\left(9.80 \mathrm{~m} / \mathrm{s}^{2}\right)\left[\sin \left(35^{\circ}\right)-(0.020) \cos \left(35^{\circ}\right)\right] \\
& =8.2 \times 10^{3} \mathrm{~N}
\end{aligned}
$$

(b) Similarly, we find that for car $\mathrm{A},\left(f_{\mathrm{R}}\right)_{\mathrm{A}}=\mu_{\mathrm{R}} m_{\mathrm{A}} g \cos \left(35^{\circ}\right)$. In the $x$-direction for car A ,

$$
\begin{aligned}
\left(F_{\text {net }}\right)_{x} & =T_{\text {motor }}+T_{\mathrm{D} \text { on A }}-\left(f_{\mathrm{R}}\right)_{\mathrm{A}}-\left(F_{\mathrm{G}}\right)_{\mathrm{A}} \sin \left(35^{\circ}\right)=0 \\
T_{\text {motor }} & =m_{\mathrm{A}} g \sin \left(35^{\circ}\right)+\mu_{\mathrm{R}} m_{\mathrm{A}} g \cos \left(35^{\circ}\right)-m_{\mathrm{D}} g \sin \left(35^{\circ}\right)+\mu_{\mathrm{R}} m_{\mathrm{D}} g \cos \left(35^{\circ}\right)
\end{aligned}
$$

Here, we have used $T_{\mathrm{A} \text { on } \mathrm{D}}=T_{\mathrm{D} \text { on } \mathrm{A}}$. If we also use $\mathrm{m}_{\mathrm{A}}=\mathrm{m}_{\mathrm{D}}$, then

$$
T_{\text {motor }}=2 \mu_{\mathrm{R}} m_{\mathrm{A}} g \cos \left(35^{\circ}\right)=4.8 \times 10^{2} \mathrm{~N}
$$

Assess: Careful examination of the free-body diagrams for cars D and A yields the observation that $T_{\text {motor }}=2\left(F_{\mathrm{R}}\right)_{\mathrm{A}}$ in order for the cars to be in dynamic equilibrium. It is a tribute to the design that the motor must only provide such a small force compared to the tension in the cable connecting the two cars.
7.46. Model: The painter and the chair are treated as a single object and represented as a particle. We assume that the rope is massless and that the pulley is massless and frictionless.

## Visualize:

## Pictorial representation



Solve: If the painter pulls down on the rope with force $F$, Newton's third law requires the rope to pull up on the painter with force $F$. This is just the tension in the rope. With our model of the rope and pulley, the same tension force $F$ also pulls up on the painter's chair. Newton's second law for (painter + chair) gives

$$
\begin{aligned}
2 F-F_{\mathrm{G}} & =\left(m_{\mathrm{P}}+m_{\mathrm{C}}\right) a \\
F & =\left(\frac{1}{2}\right)\left[\left(m_{\mathrm{P}}+m_{\mathrm{C}}\right) a+\left(m_{\mathrm{P}}+m_{\mathrm{C}}\right) g\right]=\frac{1}{2}\left(m_{\mathrm{P}}+m_{\mathrm{C}}\right)(a+g) \\
& =\left(\frac{1}{2}\right)(70 \mathrm{~kg}+10 \mathrm{~kg})\left(0.20 \mathrm{~m} / \mathrm{s}^{2}+9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=4.0 \times 10^{2} \mathrm{~N}
\end{aligned}
$$

Assess: A force of 400 N , which is approximately one-half the total gravitational force, is reasonable since the upward acceleration is small.
7.47. Model: Model Jorge as a particle and use the friction model.

Visualize:


Solve: If the Jorge pulls on the rope with force $F$, Newton's third law requires the rope to pull up on Jorge with force $F$. This is just the tension in the rope (i.e., $F=T$ ). With our model of the rope and pulley, the same tension pulls at Jorge's waist where the rope is tied. Applying Newton's second law to Jorge in the $y$-direction gives

$$
\sum(F)_{y}=n-F_{\mathrm{G}}=0 \Rightarrow n=F_{\mathrm{G}}=m g
$$

From the friction model, we have $f_{r}=\mu_{r} n=\mu_{r} m g$. Using this in the equation below, which is derived by using Newton's second law applied in the $x$-direction, gives

$$
\begin{aligned}
\sum(F)_{x} & =F+T-f_{r}=m a \\
2 F-\mu_{r} m g_{r} & =m a \\
a & =\frac{2 F}{m}-\mu_{r} g
\end{aligned}
$$

7.48. Model: Use the particle model for the tightrope walker and the rope. The rope is assumed to be massless, so the tension in the rope is uniform.

## Visualize:

## Pictorial representation



Solve: Newton's second law applied to the tightrope walker gives

$$
F_{\mathrm{R} \text { on } \mathrm{W}}-F_{\mathrm{G}}=m a \Rightarrow F_{\mathrm{R} \text { on } \mathrm{W}}=m(a+g)=(70 \mathrm{~kg})\left(8.0 \mathrm{~m} / \mathrm{s}^{2}+9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=1.25 \times 10^{3} \mathrm{~N}
$$

Newton's second law applied to the rope gives

$$
\sum\left(F_{\text {on R }}\right)_{y}=T \sin \theta+T \sin \theta-F_{\mathrm{W} \text { on R }}=0 \mathrm{~N} \Rightarrow T=\frac{F_{\mathrm{W} \text { on } \mathrm{R}}}{2 \sin \left(10^{\circ}\right)}=\frac{F_{\mathrm{R} \text { on } \mathrm{W}}}{2 \sin \left(10^{\circ}\right)}=\frac{1.25 \times 10^{3} \mathrm{~N}}{2 \sin \left(10^{\circ}\right)}=3.6 \times 10^{3} \mathrm{~N}
$$

We used $F_{\mathrm{W} \text { on } \mathrm{R}}=F_{\mathrm{R} \text { on } \mathrm{W}}$ because they are an action/reaction pair.
7.49. Model: Use the particle model for the wedge and the block.

## Visualize:

## Pictorial representation




The block will not slip relative to the wedge if they both have the same horizontal acceleration $a$. Note that $n_{1}$ on 2 and $n_{2 \text { on } 1}$ form a third law pair, so $n_{1 \text { on } 2}=n_{2 \text { on } 1}$.

Solve: Newton's second law applied to block $m_{2}$ in the $y$-direction gives

$$
\sum\left(F_{\text {on } 2}\right)_{y}=n_{1 \text { on } 2} \cos \theta-\left(F_{\mathrm{G}}\right)_{2}=0 \mathrm{~N} \Rightarrow n_{1 \text { on } 2}=\frac{m_{2} g}{\cos \theta}
$$

Combining this equation with the $x$-component of Newton's second law yields:

$$
\sum\left(F_{\text {on } 2}\right)_{x}=n_{1 \text { on } 2} \sin \theta=m_{2} a \Rightarrow a=\frac{n_{1 \text { on } 2} \sin \theta}{m_{2}}=g \tan \theta
$$

Newton's second law applied to the wedge gives

$$
\begin{gathered}
\sum\left(F_{\text {on } 1}\right)_{x}=F-n_{2} \text { on } 1 \sin \theta=m_{1} a \\
F=m_{1} a+n_{2 \text { on } 1} \sin \theta=m_{1} a+m_{2} a=\left(m_{1}+m_{2}\right) a=\left(m_{1}+m_{2}\right) g \tan \theta
\end{gathered}
$$

7.50. Model: Treat the basketball player $(\mathrm{P})$ as a particle, and use the constant-acceleration kinematic equations. Visualize:

## Pictorial representation



$$
\begin{aligned}
& \text { Known } \\
& y_{0}=v_{0 y}=t_{0}=0 \\
& y_{1}=60 \mathrm{~cm} \\
& y_{2}=140 \mathrm{~cm} \\
& m=100 \mathrm{~kg} \\
& \text { Find } \\
& \hline v_{1 v} \quad a_{0}
\end{aligned}
$$




Solve: (a) While in the process of jumping, the basketball player is pressing down on the floor as he straightens his legs. He exerts a force $F_{\mathrm{P} \text { on } \mathrm{F}}$ on the floor. The player experiences a gravitational force $\left(\vec{F}_{\mathrm{G}}\right)_{P}$ as well as a normal force from the floor $\vec{n}_{\mathrm{F} \text { on } \mathrm{P}}$. The floor experiences the force $\vec{F}_{\mathrm{P} \text { on } \mathrm{F}}$ exerted by the player.
(b) The player standing at rest exerts a force $\vec{F}_{\mathrm{P} \text { on } \mathrm{F}}$ on the floor. The normal force $\vec{n}_{\mathrm{F} \text { on } \mathrm{P}}$ is the reaction force to $\vec{F}_{\mathrm{P} \text { on } \mathrm{F}}$. But $n_{\mathrm{F} \text { on } \mathrm{P}}=F_{\mathrm{P} \text { on } \mathrm{F}}$, so $\vec{F}_{\text {net }}=0 \mathrm{~N}$. When the basketball player accelerates upward by straightening his legs, his speed has to increase from zero to $v_{1 y}$, which is the speed with which he leaves the floor. Thus, according to Newton's second law, there must be a net upward force on him during this time. This can be true only if $n_{\mathrm{F} \text { on } \mathrm{P}}>\left(F_{\mathrm{G}}\right)_{\mathrm{P}}$. In other words, the player presses on the floor with a force $F_{\mathrm{P} \text { on } \mathrm{F}}$ larger than the gravitational force on him, which is equal to his weight. The reaction force $\vec{n}_{\mathrm{F} \text { on } \mathrm{P}}$ then exceeds his weight and accelerates him upward until his feet leave the floor.
(c) The height of $80 \mathrm{~cm}=0.80 \mathrm{~m}$ is sufficient to determine the speed $v_{1 y}$ with which he leaves the floor. Once his feet are off the floor, he is simply in free fall, with $a_{1}=-g$. From kinematics,

$$
v_{2 y}^{2}=v_{1 y}^{2}+2 a_{1}\left(y_{2}-y_{1}\right) \Rightarrow 0 \mathrm{~m}^{2} / \mathrm{s}^{2}=v_{1 y}^{2}+2(-g)(0.80 \mathrm{~m}) \quad \Rightarrow \quad v_{1 y}=\sqrt{(2 g)(0.80 \mathrm{~m})}=3.96 \mathrm{~m} / \mathrm{s}
$$

The basketball player reaches $v_{1 y}=4.0 \mathrm{~m} / \mathrm{s}$ by accelerating from rest through a distance of 0.60 m .
(d) Assuming $a_{0}$ to be constant during the jump, we find

$$
v_{1 y}^{2}=v_{0 y}^{2}+2 a_{0}\left(y_{1}-y_{0}\right)=0 \mathrm{~m}^{2} / \mathrm{s}^{2}+2 a_{0}\left(y_{1}-0 \mathrm{~m}\right) \Rightarrow a_{0}=\frac{v_{1 y}^{2}}{2 y_{1}}=\frac{(3.96 \mathrm{~m} / \mathrm{s})^{2}}{2(0.60 \mathrm{~m})}=13 \mathrm{~m} / \mathrm{s}^{2}
$$

(e) The scale reads the value of $n_{\mathrm{F} \text { on } \mathrm{P}}$, which is the force exerted by the scale on the player. Before jumping,

$$
n_{\mathrm{F} \text { on } \mathrm{P}}-\left(F_{\mathrm{G}}\right)_{\mathrm{P}}=0 \mathrm{~N} \Rightarrow n_{\mathrm{F} \text { on } \mathrm{P}}=\left(F_{\mathrm{G}}\right)_{\mathrm{P}}=m g=(100 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=980 \mathrm{~N}
$$

While accelerating upward,

$$
n_{\mathrm{F} \text { on } \mathrm{P}}-m g=m a_{0} \Rightarrow n_{\mathrm{F} \text { on } \mathrm{P}}=m a_{0}+m g=m g\left(1+\frac{a_{0}}{g}\right)=(980 \mathrm{~N})\left(1+\frac{13.1}{9.8}\right)=2.3 \mathrm{kN}
$$

After leaving the scale, $n_{\text {F on } P}=0 \mathrm{~N}$ because there is no contact with the scale.
7.51. A 1.0 kg wood block is placed on top of a 2.0 kg wood block. A horizontal rope pulls the 2.0 kg block across a frictionless floor with a force of 21.0 N . Does the 1.0 kg block on top slide?

## Visualize:

## Pictorial representation



Solve: The 1.0 kg block is accelerated by static friction. It moves smoothly with the lower block if $f_{\mathrm{s}}<\left(f_{\mathrm{s}}\right)_{\max }$. It slides if the force that would be needed to keep it in place exceeds $\left(f_{\mathrm{s}}\right)_{\max }$. Begin by assuming that the blocks move together with a common acceleration $a$. Newton's second law gives

Top block: $\quad \sum\left(F_{\text {on } 1}\right)_{x}=f_{\mathrm{s}}=m_{1} a$
Bottom block: $\quad \sum\left(F_{\text {on 2 }}\right)_{x}=T_{\text {pull }}-f_{\mathrm{s}}=m_{2} a$
Adding these two equations gives $T_{\text {pull }}=\left(m_{1}+m_{2}\right) a$, or $a=(21.0 \mathrm{~N}) /(1.0 \mathrm{~kg}+2.0 \mathrm{~kg})=7.0 \mathrm{~m} / \mathrm{s}^{2}$. The static friction force needed to accelerate the top block at $7.0 \mathrm{~m} / \mathrm{s}^{2}$ is

$$
f_{\mathrm{s}} m_{1} a=(1.0 \mathrm{~kg})\left(7.0 \mathrm{~m} / \mathrm{s}^{2}\right)=7.0 \mathrm{~N}
$$

To find the maximum possible static friction force $\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} n_{1}$, the $y$-equation of Newton's second law for the top block shows that $n_{1}=m_{1} g$. Thus

$$
\left(f_{\mathrm{s}}\right)_{\max }=\mu_{\mathrm{s}} m_{1} g=(0.50)(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=4.9 \mathrm{~N}
$$

Because $7.0 \mathrm{~N}>4.9 \mathrm{~N}$, static friction is not sufficient to accelerate the top block, so it slides.
7.52. A 1.0 kg wood block is placed behind a 2.0 kg wood block on a horizontal table. The coefficients of kinetic friction with the table are 0.30 for the 1.0 kg block and 0.50 for the 2.0 kg block. The 1.0 kg block is pushed forward, against the 2.0 block, and released with a speed of $2.0 \mathrm{~m} / \mathrm{s}$. How far do the blocks travel before stopping?

## Visualize:

## Pictorial representation



Solve: The 2.0 kg block in front has a larger coefficient of friction. Thus the 1.0 kg block pushes against the rear of the 2.0 kg block and, in reaction, the 2.0 kg block pushes back against the 1.0 kg block. There's no vertical acceleration, so $n_{1}=m_{1} g$ and $n_{2}=m_{2} g$, leading to $f_{1}=\mu_{1} m_{1} g$ and $f_{2}=\mu_{2} m_{2} g$. Applying Newton's second law along the $x$-axis gives

$$
\begin{array}{ll}
1 \text { kg block: } & \sum\left(F_{\text {on } 1}\right)_{x}=-F_{2 \text { on } 1}-f_{1}=-F_{2 \text { on } 1}-\mu_{1} m_{1} g=m_{1} a \\
2 \text { kg block: } & \sum\left(F_{\text {on } 2}\right)_{x}=F_{1 \text { on } 2}-f_{2}=F_{2} \text { on } 1-\mu_{2} m_{2} g=m_{2} a
\end{array}
$$

where we used $a_{1}=a_{2}=a$. Also, $F_{1 \text { on } 2}=F_{2 \text { on } 1}$ because they are an third-law action/reaction pair. Adding these two equations gives

$$
\begin{aligned}
-\left(\mu_{1} m_{1}+\mu_{2} m_{2}\right) g & =\left(m_{1}+m_{2}\right) a \\
a & =-\frac{\mu_{1} m_{1}+\mu_{2} m_{2}}{m_{1}+m_{2}} g=-\frac{(0.30)(1.0 \mathrm{~kg})+(0.50)(2.0 \mathrm{~kg})}{1.0 \mathrm{~kg}+2.0 \mathrm{~kg}}\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=-4.25 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

We can now use constant-acceleration kinematics to find

$$
v_{1 x}^{2}=0=v_{0 x}^{2}+2 a\left(x_{1}-x_{0}\right) \Rightarrow x_{1}=-\frac{v_{0 x}^{2}}{2 a}=-\frac{(2.0 \mathrm{~m} / \mathrm{s})^{2}}{2\left(-4.25 \mathrm{~m} / \mathrm{s}^{2}\right)}=0.47 \mathrm{~m}
$$

7.53. Model: Treat the ball of clay and the block as particles.

## Visualize:

## Pictorial representation



Known

$m_{\mathrm{C}}=100 \mathrm{~g} \quad m_{\mathrm{B}}=900 \mathrm{~g} \quad t_{0}=0$
$\left(v_{\mathrm{C}}\right)_{0}=10 \mathrm{~m} / \mathrm{s} \quad\left(v_{\mathrm{B}}\right)_{0}=0 \quad t_{1}=0.010 \mathrm{~s}$
Find
$\left(v_{\mathrm{B}}\right)_{1}=\left(v_{\mathrm{C}}\right)_{1}$

Solve: (a) Forces $\vec{F}_{\mathrm{C} \text { on } \mathrm{B}}$ and $\vec{F}_{\mathrm{B} \text { on } \mathrm{C}}$ are an action/reaction pair, so $F_{\mathrm{B} \text { on } \mathrm{C}}=F_{\mathrm{C} \text { on } \mathrm{B}}$. Note that $a_{\mathrm{B}} \neq a_{\mathrm{C}}$ because the clay is decelerating while the block is accelerating. Newton's second law applied in the $x$-direction gives

$$
\begin{aligned}
& \text { Clay: } \quad \sum\left(F_{\text {on C }}\right)_{x}=-F_{\mathrm{B} \text { on } \mathrm{C}}=m_{\mathrm{C}} a_{\mathrm{C}} \\
& \text { Block: } \quad \sum\left(F_{\text {on B }}\right)_{x}=F_{\mathrm{C} \text { on } \mathrm{B}}=F_{\mathrm{B} \text { on } \mathrm{C}}=m_{\mathrm{B}} a_{\mathrm{B}}
\end{aligned}
$$

Equating the two expressions for $F_{\mathrm{B} \text { on } \mathrm{C}}$ gives

$$
a_{\mathrm{C}}=-\frac{m_{\mathrm{B}}}{m_{\mathrm{C}}} a_{\mathrm{B}}
$$

Turning to kinematics, the velocity of each after $\Delta t$ is

$$
\begin{gathered}
\left(v_{\mathrm{C}}\right)_{1}=\left(v_{\mathrm{C}}\right)_{0}+a_{\mathrm{C}} \Delta t \\
\left(v_{\mathrm{B}}\right)_{1}=\left(v_{\mathrm{B}}\right)_{0}+a_{\mathrm{B}} \Delta t=a_{\mathrm{B}} \Delta t
\end{gathered}
$$

But $\left(v_{\mathrm{C}}\right)_{1}=\left(v_{\mathrm{B}}\right)_{1}$ because the clay and the block are moving together after $\Delta t$ has elapsed. Equating these two expressions gives $\left(v_{\mathrm{C}}\right)_{0}+a_{\mathrm{C}} \Delta t=a_{\mathrm{B}} \Delta t$, from which we find

$$
a_{\mathrm{C}}=a_{\mathrm{B}}-\frac{\left(v_{\mathrm{C}}\right)_{0}}{\Delta t}
$$

We can now equate the two expressions for $a_{\mathrm{C}}$ :

$$
-\frac{m_{\mathrm{B}}}{m_{\mathrm{C}}} a_{\mathrm{B}}=a_{\mathrm{B}}-\frac{\left(v_{\mathrm{C}}\right)_{0}}{\Delta t} \Rightarrow a_{\mathrm{B}}=\frac{\left(v_{\mathrm{C}}\right)_{0} / \Delta t}{1+m_{\mathrm{B}} / m_{\mathrm{C}}}=\frac{(10 \mathrm{~m} / \mathrm{s}) /(0.010 \mathrm{~s})}{1+(900 \mathrm{~g}) / 100 \mathrm{~g}}=100 \mathrm{~m} / \mathrm{s}^{2}
$$

Then $a_{\mathrm{C}}=-9 a_{\mathrm{B}}=-900 \mathrm{~m} / \mathrm{s}^{2}$. With the acceleration now known, we can use either kinematic equation to find

$$
\left(v_{\mathrm{C}}\right)_{1}=\left(v_{\mathrm{B}}\right)_{1}=\left(100 \mathrm{~m} / \mathrm{s}^{2}\right)(0.010 \mathrm{~s})=1.0 \mathrm{~m} / \mathrm{s}
$$

(b) $F_{\mathrm{C} \text { on } \mathrm{B}}=m_{\mathrm{B}} a_{\mathrm{B}}=(0.90 \mathrm{~kg})\left(100 \mathrm{~m} / \mathrm{s}^{2}\right)=90 \mathrm{~N}$.
(c) $F_{\mathrm{B} \text { on } \mathrm{C}}=m_{\mathrm{C}} a_{\mathrm{C}}=(0.10 \mathrm{~kg})\left(900 \mathrm{~m} / \mathrm{s}^{2}\right)=90 \mathrm{~N}$.

Assess: The two forces are of equal magnitude, as expected from Newton's third law.
7.54. Model: Use the particle model for the two blocks. Assume a massless rope and massless, frictionless pulleys. Visualize:

Pictorial representation


Note that for every meter block 1 moves forward, one meter is provided to block 2 . So each rope on $m_{2}$ has to be lengthened by one-half meter. Thus, the acceleration constraint is $a_{2}=-\frac{1}{2} a_{1}$.
Solve: Newton's second law applied to block 1 gives $T=m_{1} a_{1}$. Newton's second law applied to block 2 gives $2 T-\left(F_{\mathrm{G}}\right)_{2}=m_{2} a_{2}$. Combining these two equations gives

$$
2\left(m_{1} a_{1}\right)-m_{2} g=m_{2}\left(-\frac{1}{2} a_{1}\right) \Rightarrow a_{1}\left(4 m_{1}+m_{2}\right)=2 m_{2} g \quad \Rightarrow \quad a_{1}=\frac{2 m_{2} g}{4 m_{1}+m_{2}}
$$

where we have used $a_{2}=-\frac{1}{2} a_{1}$.
Assess: If $m_{1}=0 \mathrm{~kg}$, then $a_{2}=-g$. This is what is expected for a freely falling object.
7.55. Model: Use the particle model for the two blocks. Assume a massless rope and massless, frictionless pulleys. Visualize:

Pictorial representation


For every one meter that the $1.0-\mathrm{kg}$ block goes down, each rope on the $2.0-\mathrm{kg}$ block will be shortened by one-half meter. Thus, the acceleration constraint is $a_{1}=-2 a_{2}$.
Solve: Newton's second law applied to the two blocks gives

$$
2 T=m_{2} a_{2} T-\left(F_{\mathrm{G}}\right)_{1}=m_{1} a_{1}
$$

Since $a_{1}=-2 a_{2}$, the above equations become

$$
\begin{aligned}
2 T & =m_{2} a_{2} T-m_{1} g=m_{1}\left(-2 a_{2}\right) \\
m_{2} \frac{a_{2}}{2}+m_{1}\left(2 a_{2}\right) & =m_{1} g \\
a_{2} & =\frac{2 m_{1} g}{m_{2}+4 m_{1}}=\frac{2(1.0 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)}{(2.0 \mathrm{~kg}+4.0 \mathrm{~kg})}=3.3 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

Assess: If $m_{1}=0 \mathrm{~kg}$, then $a_{2}=0 \mathrm{~m} / \mathrm{s}^{2}$, which is expected.
7.56. Model: The hamster of mass $m$ and the wedge with mass $M$ will be treated as objects 1 and 2, respectively. They will be treated as particles.

## Visualize:

Pictorial representation

Known
$m=200 \mathrm{~g}$
$M=800 \mathrm{~g}$
$\theta=40^{\circ}$
Find

| $n_{2}$, when hamster does |
| :--- |
| not move and then when |
| he slides down. |



The scale is denoted by the letter s.
Solve: (a) The reading of the scale is the magnitude of the force $\vec{n}_{2}$ that the scale exerts upward. There are two action/reaction pairs. Initially the hamster of mass $m$ is stuck in place and is in static equilibrium with $\vec{F}_{\text {net }}=0 \mathrm{~N}$. Because of the shape of the blocks, it is not clear whether the scale has to exert a horizontal friction force $\vec{f}_{\text {s on } 2}$ to prevent horizontal motion. We've included one just in case. Newton's second law for the hamster is

$$
\begin{aligned}
\left(F_{\text {net on } 1}\right)_{x}=m g \sin \theta-f_{2 \text { on } 1}=0 \mathrm{~N} & \Rightarrow f_{2 \text { on } 1}=m g \sin \theta \\
\left(F_{\text {net on } 1}\right)_{y}=n_{1}-m g \cos \theta=0 \mathrm{~N} & \Rightarrow n_{1}=m g \cos \theta
\end{aligned}
$$

For the wedge, we see from Newton's third law that $n_{1}^{\prime}=n_{1}=m g \cos \theta$ and that $f_{2 \text { on } 1}=f_{1 \text { on } 2}=m g \sin \theta$. Using these equations, Newton's second law for the wedge is

$$
\begin{gathered}
\left(F_{\text {net on } 2)_{x}}=f_{1 \text { on } 2} \cos \theta+f_{\text {s on } 2}-n_{1}^{\prime} \sin \theta=m g \sin \theta \cos \theta+f_{\text {s on } 2}-m g \cos \theta \sin \theta=0 \mathrm{~N} \Rightarrow f_{\text {s on } 2}=0 \mathrm{~N}\right. \\
\\
\left(F_{\text {net on } 2}\right)_{y}=n_{2}-n_{1}^{\prime} \cos \theta-f_{1 \text { on } 2} \sin \theta-M g=n_{2}-m g \cos ^{2} \theta-m g \sin ^{2} \theta-M g=0 \mathrm{~N} \\
\\
n_{2}=m g\left(\cos ^{2} \theta+\sin ^{2} \theta\right)+M g=(M+m) g=(0.800 \mathrm{~kg}+0.200 \mathrm{~kg})\left(9.8 \mathrm{~m} / \mathrm{s}^{2}\right)=9.8 \mathrm{~N}
\end{gathered}
$$

First we find that $f_{\text {son } 2}=0 \mathrm{~N}$, so no horizontal static friction is needed to prevent motion. More interesting, the scale reading is $(M+m) g$ which is the total gravitational force resting on the scale. This is the expected result.
(b) Now suppose that the hamster is accelerating down the wedge. The total mass is still $M+m$, but is the reading still $(M+m) g$ ? The frictional forces between the systems 1 and 2 have now vanished, and system 1 now has an acceleration. However, the acceleration is along the hamster's $x$-axis, so $a_{1 y}=0 \mathrm{~m} / \mathrm{s}^{2}$. The hamster's $y$-equation is still

$$
\left(F_{\text {net on } 1}\right)_{y}=n_{1}-m g \cos \theta=0 \mathrm{~N} \Rightarrow n_{1}=m g \cos \theta
$$

We still have $n_{1}^{\prime}=n_{1}=m g \cos \theta$, so the $y$-equation for block 2 (with $a_{2 y}=0 \mathrm{~m} / \mathrm{s}^{2}$ ) is

$$
\begin{aligned}
\left(F_{\text {net on } 2}\right)_{y} & =n_{2}-n_{1}^{\prime} \cos \theta-M g=n_{2}-m g \cos ^{2} \theta-M g=0 \mathrm{~N} \\
n_{2} & =m g \cos ^{2} \theta+M g=\left(M+m \cos ^{2} \theta\right) g=9.0 \mathrm{~N}
\end{aligned}
$$

Assess: The scale reads less than it did when the hamster was at rest. This makes sense if you consider the limit $\theta \rightarrow 90^{\circ}$, in which case $\cos \theta \rightarrow 0$. If the face of the wedge is vertical, then the hamster is simply in free fall and can have no effect on the scale (at least until impact!). So for $\theta=90^{\circ}$ we expect the scale to record $M g$ only, and that is indeed what the expression for $n_{2}$ gives.
7.57. Model: The hanging masses $m_{1}, m_{2}$, and $m_{3}$ are modeled as particles. Pulleys A and B are massless and frictionless. The strings are massless.

## Visualize:

Pictorial representation


Solve: (a) The length of the string over pulley $B$ is constant. Therefore,

$$
\left(y_{\mathrm{B}}-y_{3}\right)+\left(y_{\mathrm{B}}-y_{\mathrm{A}}\right)=L_{\mathrm{B}} \quad \Rightarrow \quad y_{\mathrm{A}}=2 y_{\mathrm{B}}-y_{3}-L_{\mathrm{B}}
$$

The length of the string over pulley A is constant. Thus,

$$
\begin{aligned}
\left(y_{\mathrm{A}}-y_{2}\right)+\left(y_{\mathrm{A}}-y_{1}\right) & =L_{\mathrm{A}}=2 y_{\mathrm{A}}-y_{1}-y_{2} \\
2\left(2 y_{\mathrm{B}}-y_{3}-L_{\mathrm{B}}\right)-y_{1}-y_{2} & =L_{\mathrm{A}} \quad \Rightarrow 2 y_{3}+y_{2}+y_{1}=\mathrm{constant}
\end{aligned}
$$

This constraint implies that

$$
2 \frac{d y_{3}}{d t}+\frac{d y_{2}}{d t}+\frac{d y_{1}}{d t}=0 \mathrm{~m} / \mathrm{s}=2 v_{3 y}+v_{2 y}+v_{1 y}
$$

Also by differentiation, $2 a_{3 y}+a_{2 y}+a_{1 y}=0 \mathrm{~m} / \mathrm{s}^{2}$.
(b) Applying Newton's second law to the masses $m_{3}, m_{2}, m_{1}$, and pulley A gives

$$
T_{\mathrm{B}}-m_{3} g=m_{3} a_{3 y} \quad T_{\mathrm{A}}-m_{2} g=m_{2} a_{2 y} \quad T_{\mathrm{A}}-m_{1} g=m_{1} a_{1} y \quad T_{\mathrm{B}}-2 T_{\mathrm{A}}=0 \mathrm{~N}
$$

The pulley equation is zero because the pulley is massless. These four equations plus the acceleration constraint consitute five equations with five unknowns (two tensions and three accelerations). To solve for $T_{\mathrm{A}}$, multiply the $m_{3}$ equation by 2 , substitute $2 T_{\mathrm{B}}=4 T_{\mathrm{A}}$, then divide each of the mass equations by the mass. This gives the three equations

$$
\begin{aligned}
4 T_{\mathrm{A}} / m_{3}-2 g & =2 a_{3 y} \\
T_{\mathrm{A}} / m_{2}-g & =a_{2 y} \\
T_{\mathrm{A}} / m_{1}-g & =a_{1 y}
\end{aligned}
$$

If these three equations are added, the right side adds to zero because of the acceleration constraint. Thus

$$
\left(4 / m_{3}+1 / m_{2}+1 / m_{2}\right) T_{\mathrm{A}}-4 g=0 \quad \Rightarrow \quad T_{\mathrm{A}}=\frac{4 g}{\left(4 / m_{3}+1 / m_{2}+1 / m_{2}\right)}
$$

(c) Using numerical values, we find $T_{\mathrm{A}}=18.97 \mathrm{~N}$. Then

$$
\begin{aligned}
& a_{1 y}=T_{\mathrm{A}} / m_{1}-g=-2.2 \mathrm{~m} / \mathrm{s}^{2} \\
& a_{2 y}=T_{\mathrm{A}} / m_{2}-g=2.9 \mathrm{~m} / \mathrm{s}^{2} \\
& a_{3 y}=2 T_{\mathrm{A}} / m_{3}-g=-0.32 \mathrm{~m} / \mathrm{s}^{2}
\end{aligned}
$$

(d) $m_{3}=m_{1}+m_{2}$, so it appears at first as if $m_{3}$ should hang in equilibrium. For this to be so, tension $T_{\mathrm{B}}$ would need to equal $m_{3} g$. However, $T_{\mathrm{B}}$ is not $\left(m_{1}+m_{2}\right) g$ because masses $m_{1}$ and $m_{2}$ are accelerating rather than hanging at rest. Consequently, tension $T_{\mathrm{B}}$ is not able to balance the weight of $m_{3}$.


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